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# Modeling lightning- $\text{NO}_x$ chemistry on a sub-grid scale in a global chemical transport model

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**Abstract.** For the first time, a plume-in-grid approach is implemented in a chemical transport model (CTM) to parameterize the effects of the nonlinear reactions occurring within high concentrated  $\text{NO}_x$  plumes from lightning  $\text{NO}_x$  emissions ( $\text{LNO}_x$ ) in the upper troposphere. It is characterized by a set of parameters including the plume lifetime, the effective reaction rate constant related to  $\text{NO}_x$ – $\text{O}_3$  chemical interactions, and the fractions of  $\text{NO}_x$  conversion into  $\text{HNO}_3$  within the plume. Parameter estimates were made using the Dynamical Simple Model of Atmospheric Chemical Complexity (DSMACC) box model, simple plume dispersion simulations, and the 3-D Meso-NH (non-hydrostatic mesoscale atmospheric model). In order to assess the impact of the  $\text{LNO}_x$  plume approach on the  $\text{NO}_x$  and  $\text{O}_3$  distributions on a large scale, simulations for the year 2006 were performed using the GEOS-Chem global model with a horizontal resolution of  $2^\circ \times 2.5^\circ$ . The implementation of the  $\text{LNO}_x$  parameterization implies an  $\text{NO}_x$  and  $\text{O}_3$  decrease on a large scale over the region characterized by a strong lightning activity (up to 25 and 8 %, respectively, over central Africa in July) and a relative increase downwind of  $\text{LNO}_x$  emissions (up to 18 and 2 % for  $\text{NO}_x$  and  $\text{O}_3$ , respectively, in July). The calculated variability in  $\text{NO}_x$  and  $\text{O}_3$  mixing ratios around the mean value according to the known uncertainties in the parameter estimates is at a maximum over continental tropical regions with  $\Delta\text{NO}_x$  [−33.1, +29.7] ppt and  $\Delta\text{O}_3$  [−1.56, +2.16] ppb, in January, and  $\Delta\text{NO}_x$  [−14.3, +21] ppt and  $\Delta\text{O}_3$  [−1.18, +1.93] ppb, in July, mainly depending on the determination of the diffusion properties of the atmosphere

and the initial  $\text{NO}$  mixing ratio injected by lightning. This approach allows us (i) to reproduce a more realistic lightning  $\text{NO}_x$  chemistry leading to better  $\text{NO}_x$  and  $\text{O}_3$  distributions on the large scale and (ii) to focus on other improvements to reduce remaining uncertainties from processes related to  $\text{NO}_x$  chemistry in CTM.

## 1 Introduction

Lightning emissions are one of the most important sources of nitrogen oxides ( $\text{NO}_x \equiv \text{NO} + \text{NO}_2$ ) in the upper troposphere (WMO, 1999; Hudman et al., 2007). Lightning primarily produces  $\text{NO}$  and may also induce a negligible quantity of  $\text{NO}_2$  with a ratio  $\text{NO}_2/\text{NO}_x$  of 0.5 to 0.1 (Franzblau, 1991; Stark et al., 1996).  $\text{NO}_x$  emitted by lightning ( $\text{LNO}_x$ ) impacts the tropospheric ozone burden (Stockwell et al., 1999; Hauglustaine et al., 2001; Grewe, 2007) and the hydroxyl-radical ( $\text{OH}$ ) concentrations influencing the oxidizing capacity of the atmosphere (Labrador et al., 2004; Banerjee et al., 2014). Most  $\text{NO}_x$  produced by lightning is detrained into the free and upper troposphere, where ozone production efficiencies (OPEs) per unit  $\text{NO}_x$  emitted are 4 to 20 times higher than at the surface (Sauvage et al., 2007a; Martin et al., 2007), and therefore lightning exerts a disproportionately stronger effect on photochemistry than surface emissions (Pickering et al., 1990; Hauglustaine et al., 1994; Zhang et al., 2003; Choi et al., 2009). The longer  $\text{NO}_x$  lifetime in the upper troposphere (1–2 weeks) allows

the long-range transport of LNO<sub>x</sub> through large circulation patterns (Hemispheric Transport of Air Pollution, HTAP, report, 2010: <http://www.htap.org/>).

Although the importance of the LNO<sub>x</sub> emissions for the upper tropospheric chemistry is well known, it remains highly uncertain, with a best estimate of  $2\text{--}8\text{ TgN} \times \text{yr}^{-1}$  (Schumann and Huntrieser, 2007). Lightning NO<sub>x</sub> emissions are associated with deep convection (horizontal scale  $\sim 10\text{ km}$ ) and correspond to the “sub-grid” in global chemical transport models (horizontal resolution  $\sim 100\text{ km}$ ). Therefore, lightning NO<sub>x</sub> production must be parameterized for inclusion into a large-scale model. Global models commonly use convection proxies such as the cloud top height (Price and Rind, 1992) and the updraft intensity to estimate the lightning flashes. Flashes simulated by chemical transport models (CTMs) are commonly constrained by satellite observations (Sauvage et al., 2007b; Murray et al., 2012) from the spaceborne Lightning Imaging Sensor (LIS) in the Tropical Rainfall Measuring Mission (TRMM) and the Optical Transient Detector (OTD) (Christian et al., 2003; Tost et al., 2007). The lightning NO<sub>x</sub> emissions are then redistributed according to a vertical profile – generally a reverse “C-Shape” profile (Ott et al., 2010) – a priori defined depending on season, latitude, and continent/ocean location. Also, corrections to the calculations of lightning NO<sub>x</sub> emissions using satellite observations (SCIAMACHY, SCanning Imaging Absorption SpectroMeter for Atmospheric CHarotographY; Martin et al., 2007) and in situ measurements (INTEX-NA, Intercontinental Chemical Transport Experiment – North America; Hudman et al., 2007) are usually applied.

Despite the necessity of including lightning NO<sub>x</sub> emissions in global models, the small-scale nature of the flashes and the nonlinear chemistry (Lin et al., 1988) of the atmosphere will lead to biases on the large scale with instantaneous dilution of gases in the large grid box volume. It seems likely that this will lead to an overestimate of the OPE and an underestimate of the nitric acid (HNO<sub>3</sub>) production. For instance, by forcing NO<sub>x</sub> concentration in a GEOS-Chem grid box over southeast Asia to represent the measured lightning plumes, Cooper et al. (2014) estimated a ratio for O<sub>3</sub> to HNO<sub>3</sub> produced leading to a  $15\text{ mol mol}^{-1}$  OPE in lightning plumes, which reinforces the fact that instantaneous dilution in global model implies issues in sub-grid chemistry.

In this work, a more realistic lightning NO<sub>x</sub> chemistry as well as a plume parameterization is implemented into a global CTM allowing us to reproduce the NO<sub>x</sub> and O<sub>3</sub> distributions more accurately on a large scale. The plume approach used in this study was previously developed by Cariolle et al. (2009) for aircraft NO<sub>x</sub> emissions in the LMDz-INCA (Laboratoire de Météorologie Dynamique, Interaction avec la Chimie et les Aérosols) and MOBIDIC (MOBile DIgital Compute) models and also implemented to deal with ship NO<sub>x</sub> emissions (Huszar et al., 2010). This approach avoids the double count in the CTM calculation of

the emitted NO<sub>x</sub>, which is first instantaneously diluted onto the point grid and secondly as the plume form. In addition, the plume parameterization is the first that considers the NO<sub>x</sub> from lightning as a plume, with the transport of the related nonlinear chemistry effects. NO<sub>x</sub> from lightning emissions are emitted in the upper troposphere, characterized by strong winds allowing the large-scale transport of trace species. Thus, it is relevant to consider plume growth from lightning emissions, which may have been diluted a long time after the initial lightning pulse, downwind of emissions. Consequently, the plume parameterization previously developed for aircraft exhausts has been adjusted to LNO<sub>x</sub> emissions and implemented into the GEOS-Chem global chemical transport model.

Section 2 gives a description of the GEOS-Chem model in which the plume-in-grid parameterization is implemented and the models which are used to evaluate the diffusion properties of the atmosphere and to determine parameters characterizing the physics and chemistry of the lightning NO<sub>x</sub> plume. A concise description of the plume approach is then presented in Sect. 3, followed by a detailed explanation of the determination of parameters related to LNO<sub>x</sub> emissions. Section 4 summarizes the results of the simulations performed with GEOS-Chem and finally these results and the sensitivity to NO<sub>x</sub> and O<sub>3</sub> variations of the parameterization are discussed in Sect. 5.

## 2 Models

Three different models are used in this evaluation and are described in this section. GEOS-Chem is used to provide a global framework to assess the impact of lightning NO<sub>x</sub>. Meso-NH (non-hydrostatic mesoscale atmospheric model) is used to provide estimates of the plume diffusion timescales, and Dynamical Simple Model of Atmospheric Chemical Complexity (DSMACC) is a box model used to assess the nonlinear chemistry in the plume.

### 2.1 The GEOS-Chem chemical transport model

The GEOS-Chem chemical transport model (Bey et al., 2001) is a global 3-D model of atmospheric composition driven by an assimilated meteorology from the Goddard Earth Observing System (GEOS-5) of the NASA Global Modeling Assimilation Office (GMAO). The 09-01-01 version ([http://wiki.seas.harvard.edu/geos-chem/index.php/GEOS-Chem\\_v9-01-01](http://wiki.seas.harvard.edu/geos-chem/index.php/GEOS-Chem_v9-01-01)) of the CTM has been used in this study. The model transports 43 tracers to describe tropospheric O<sub>3</sub>–NO<sub>x</sub>–VOC (volatile organic compound) chemistry. The horizontal resolution is  $2^\circ \times 2.5^\circ$ , and 47 vertical levels are defined from the ground to 80 km altitude. The CTM includes modules for emissions, transport, chemistry, deposition, aerosols, and surface.

The large-scale advection of tracers is performed using the TPCORE advection scheme (Lin and Rood, 1996), corresponding to a semi-Lagrangian flux method. Shallow and deep moist convection processes are carried out using the relaxed Arakawa-Schubert scheme (Moorthi and Suarez, 1991). Mixing in the lower atmospheric layers is represented by a nonlocal scheme of the planetary boundary layer described by Lin and McElroy (2010). The wet deposition for water-soluble aerosols and for gases follows Liu et al. (2001) and Amos et al. (2012). Aerosol scavenging by ice crystals and cold or mixed precipitation is also reproduced in the model (Wang et al., 2011). The dry deposition is associated with a scheme which calculates bulk surface resistance in series (Wesely, 1989). Photolysis rates are calculated with the Fast-JX code (Bian and Prather, 2002). The atmospheric chemistry is resolved using the SMVGEAR (sparse-matrix vectorized gear code) solver (Jacobson and Turco, 1994), with more than 300 species and 785 chemical reactions. Heterogeneous chemical reactions are represented on the surface of aerosols (Bey et al., 2001; Martin et al., 2002). The effects of aerosols on the photolysis rates are based on Martin et al. (2003). Primary NO<sub>x</sub> and VOC emissions are separated depending on sources. Global anthropogenic emissions are given by the GEIA (Global Emissions Initiative; Wang et al., 1998) and EGDAR (Emission Database for Global Atmospheric Research; Olivier, 2005) inventories, and regional anthropogenic emissions exceed those for the US (National Emissions Inventory, NEI05), Canada (Criteria Air Contaminants, CAC), Mexico (Big Bend Regional Aerosol and Visibility Observational study, BRAVO), Europe (EMEP, European Monitoring and Evaluation Programme), and east Asia (Streets et al., 2006; Zhang et al., 2009). Biofuel emissions are provided by the EPA (Environmental Protection Agency) and STREETS 2006 inventories (Yevich and Logan, 2003), biomass burning emissions by the GFED (Global Fire Emissions Database) inventory (van der Werf et al., 2010), and biogenic emissions by the MEGAN (Model of Emissions of Gases and Aerosols from Nature) model calculations (Guenther et al., 2012). In addition, NO<sub>x</sub> from soil emissions is calculated by an algorithm depending on temperature and precipitation (Yienger and Levy, 1995).

In order to calculate the NO<sub>x</sub> from lightning, flash rates are first calculated in active deep convection using the Price and Rind scheme based on cloud top height (Price and Rind, 1992, 1994); then flash rates are adjusted with local scaling factors to match the satellite climatology (Sauvage et al., 2007b; Murray et al., 2012), and the total column emissions are determined using NO<sub>x</sub> yields that differ in the tropics and northern extratropics. Finally, the total column is distributed vertically using the reverse C-shaped profile from Ott et al. (2010). Note that the base lightning NO<sub>x</sub> scheme is described in detail by Murray et al. (2012).

## 2.2 The Meso-NH model

The Meso-NH model is an atmospheric model developed jointly by the Laboratoire d'Aérodynamique and by CNRM-GAME (<http://mesonh.aero.obs-mip.fr/mesonh51>). The model includes a non-hydrostatic and anelastic system of equations (Lafore et al., 1998) and has a complete set of parameterizations allowing us to reproduce physical processes such as radiation (Gregory et al., 2000), atmospheric turbulence (Cuxart et al., 1999), convection (Bechtold et al., 2000), microphysics related to warm clouds (Cohard and Pinty, 2000), and atmospheric ice (Pinty and Jabouille, 1999; Lascaux et al., 2006). Meso-NH includes also online chemistry (Tulet et al., 2003, 2006). The model deals with large (synoptic) to small (large eddy) scales. In this study, the Mesonh-49 version was used in order to compare the horizontal diffusion coefficient ( $D_h$ ) estimate within the anvil of thunderstorms from in situ measurements to a modeling ideal case of a convective cell.

## 2.3 The DSMACC box model

The DSMACC is a simple box model developed for improving our understanding of the tropospheric chemistry (Emmerson and Evans, 2009). The model is composed of the KPP (kinetic preprocessor) chemical preprocessor (Damian et al., 2002) to solve differential equations representing the chemical system. The TUV (Tropospheric Ultraviolet and Visible Radiation Model) photolysis scheme is used, which calculates the spectral irradiance, the spectral actinic flux, photodissociation coefficients ( $J$  values) (Madronich and Flocke, 1999), and biologically effective irradiance. The chemical scheme used derives from the Master Chemical Mechanism (MCM, <http://mcm.leeds.ac.uk/MCM/>), (Jenkin et al., 1997; Saunders et al., 2003), which contains 17 000 elementary reactions of 6700 primary, secondary, and radical species.

In order to study the chemical interactions that can occur in the undiluted plume fraction, a set of short simulations was carried out with the DSMACC box model as explained in Sect. 3.2.2.

## 2.4 The simple plume dispersion model

To model the dispersion of lightning NO<sub>x</sub> emissions we use a simple dispersion model similar to the plume model used for aircraft NO<sub>x</sub> emissions, except that the plume is supposed to be oriented along a vertical axis. The plume is represented as a cylinder that encompasses horizontal diffusion with a constant coefficient  $D_h$  (Sect. 3.2.1). This simple model is composed of 30 horizontal circles with spacing increasing progressively from the center axis. The discretization of the diffusion equation is mass conservative.

The chemistry scheme and associated reaction rate constants are adapted from the large-scale chemical model

MOCAGE (Modèle de Chimie Atmosphérique de Grande Echelle; Teyssède et al., 2007). It includes the main reactions involved in the NO<sub>x</sub>–HO<sub>x</sub> system. Simple plume simulations were performed in order to estimate the physical and chemical characteristics of the plumes related to lightning NO<sub>x</sub> emissions.

### 3 Plume parameterization for lightning NO<sub>x</sub> emissions

#### 3.1 General description

The LNO<sub>x</sub> plume parameterization is based on a method initially developed by Cariolle et al. (2009) for NO<sub>x</sub> emissions related to aircraft exhausts later adapted to ship emissions of NO<sub>x</sub> (Huszar et al., 2010). In this approach, the plume effects on a sub-grid scale are represented via a fuel tracer in order to follow the amount of the emitted species in the plume and an effective reaction rate for the ozone production and nitric acid production and destruction during the plume's dilution in the background (Cariolle et al., 2009; Paoli et al., 2011). The parameterization requires a proper estimation of the characteristic plume lifetime, during which the nonlinear interactions between species are important and simulated via specific rates of conversion. The approach ensures the mass conservation of species in the model. This is the only method which considers a plume evolution related to the local NO<sub>x</sub> emissions, allowing the transport of the nonlinear effects which occur on a smaller scale than the model grid.

##### 3.1.1 Physical plume formulation

Following Cariolle et al. (2009), a passive tracer (from the perspective of the usual model chemistry) is added to the CTM to represent NO<sub>x</sub> emitted by lightning. The LNO<sub>x</sub> tracer initial mass corresponds to the NO<sub>x</sub> mass at the start time of the simulation. Rather than increasing the concentration of NO<sub>x</sub> within the CTM, lightning NO<sub>x</sub> emissions now increase the concentration of this new passive tracer, which is transported in the standard way by advection and turbulence. Plume chemistry is considered to be significant when the mixing ratio of the lightning NO<sub>x</sub> tracer is higher than the critical NO<sub>x</sub> content, hereafter denoted  $r_1$ . Above this value the lightning NO<sub>x</sub> tracer is transferred to the normal NO<sub>x</sub> tracer at a rate described by a plume lifetime ( $\tau$ ), which is an exponential decay constant. This corresponds to an exchange timescale between the lightning NO<sub>x</sub> plume and the background NO<sub>x</sub>. The continuity equation related to the tracer evolution is detailed by Eq. (1).

$$\frac{\partial \overline{r_{\text{LNO}_x}}}{\partial t} + \langle F_{\text{LNO}_x} \rangle = I - \frac{1}{\tau} \cdot \overline{r_{\text{LNO}_x}}, \quad (1)$$

where  $\overline{r_{\text{LNO}_x}}$  is the mixing ratio (in ppb) of the NO<sub>x</sub> lightning tracer in the model grid (note that all overlined terms refer to grid average quantities in the CTM),  $F_{\text{LNO}_x} \equiv \nabla \times$

$(\overline{r_{\text{LNO}_x}} u) + \nabla \times (D_t \nabla \overline{r_{\text{LNO}_x}})$  and corresponds to the flux divergence related to the large-scale transport of the tracer (advection and turbulent diffusion, in molecules cm<sup>−2</sup> s<sup>−1</sup>),  $I$  is the injection rate of LNO<sub>x</sub> (in s<sup>−1</sup>), and  $\tau$  is the plume lifetime (in seconds).

The calculation of  $\tau$  requires evaluating the mass fraction of the lightning NO<sub>x</sub> ( $M(t)$ ) corresponding to the undiluted fraction of the plume and characterized by an NO<sub>x</sub> mixing ratio above the  $r_1$  critical value. In other words, the plume boundary is defined by the critical value  $r_1$  depending on the time of day. The NO<sub>x</sub> mass,  $M(t)$ , decreases monotonically to zero until  $t = T_1$  for which the tracer mixing ratio is everywhere below the  $r_1$  threshold. The plume lifetime is obtained by an exponential function depending on the mass (Eqs. 2 and 3):

$$M(t) = \int_{V_p} \rho \cdot r_p \cdot dV, \quad (2)$$

$$\tau = \int_{t_0=0}^{+\infty} \exp(-t/\tau) \cdot dt = \frac{1}{M(t_0)} \int_{t_0=0}^{T_1} M(t) \cdot dt. \quad (3)$$

Where  $V_p$  is the volume of the plume,  $\rho$  is the density of the air,  $r_p$  is the NO<sub>x</sub> mixing ratio within the plume (in ppb), and  $T_1$  is the time for which the mixing ratio  $r_p$  is below the critical value  $r_1$  everywhere. The calculation of the plume lifetime, by simple plume dispersion simulations, depends on (i) the initial emissions of NO<sub>x</sub> by lightning, (ii) the  $r_1$  value, and (iii) the dispersion properties of the atmosphere (related to the horizontal diffusion coefficient,  $D_h$ ) and is detailed in Sect. 3.2.3. Note that the mean dispersion properties of the atmosphere were associated with the horizontal diffusion only. The lightning NO<sub>x</sub> emissions occur in the convective part of clouds where the vertical diffusion is strong. Therefore, the vertical diffusion coefficient is a determining parameter for the LNO<sub>x</sub> distribution in the cloud. As mentioned in Sect. 2.1, the vertical distribution of LNO<sub>x</sub> is calculated a priori from Ott et al. (2010) as a reverse C-shaped profile. The LNO<sub>x</sub> plume parameterization is applied a posteriori; after that, lightning NO<sub>x</sub> is vertically prescribed and concerns convective outflow where NO<sub>x</sub> is detrained in the troposphere. In this region of detrainment, the horizontal dispersion may be more efficient than the vertical one as discussed in Cariolle et al. (2009).

##### 3.1.2 Plume chemistry of NO<sub>x</sub>, O<sub>3</sub>, and HNO<sub>3</sub>

Once the lightning NO<sub>x</sub> is emitted, it is transferred to model's background NO<sub>x</sub> based on the lifetime of the plume ( $\tau$ ). Thus, the continuity equation for the NO<sub>x</sub> species emitted in the plume and released to the large scale can be deduced as described by Eq. (4).

$$\frac{\partial \overline{r_{\text{NO}_x}}}{\partial t} + \langle F_{\text{NO}_x} \rangle = + \frac{1}{\tau} \cdot \overline{r_{\text{LNO}_x}} \cdot \alpha_{\text{NO}_x} + L_{\text{ss}}, \quad (4)$$

where  $\overline{r_{\text{NO}_x}}$  is the mixing ratio of NO<sub>x</sub> (in ppb) in the model grid,  $\alpha_{\text{NO}_x}$  is the molecular mass ratio between the air and NO<sub>x</sub> species, and  $L_{\text{ss}}$  is the large-scale sources and sinks (in molecules cm<sup>-2</sup> s<sup>-1</sup>), such as natural and anthropogenic emissions, photochemical reaction, mixing, and conversion to reservoir species.

We consider a fairly simple chemistry within the plume as described below. The increase in the nitrogen oxide concentration in the upper troposphere leads to ozone production through the reaction of NO with peroxide (HO<sub>2</sub>), CH<sub>3</sub>O<sub>2</sub>, or RO<sub>2</sub> radicals from the OH oxidation as shown by Reaction (R1).



In the case of large NO<sub>x</sub> injection by lightning, the NO<sub>x</sub> content ( $\sim 40$  ppt in unpolluted atmosphere) becomes close (a few ppb, according to in situ measurements; Dye et al., 2000; Huntrieser et al., 2002) to the surrounding ozone ( $60 \pm 24$  ppb) (Jaéglé et al., 1998). The ozone evolution within the plume is described by Reactions (R2)–(R6).



From these equations we can define an O<sub>x</sub> family (O<sub>x</sub>  $\equiv$  O + O<sub>3</sub> + NO<sub>2</sub>) where the only net loss of O<sub>x</sub> is by reactions between atomic oxygen and NO<sub>2</sub> or O<sub>3</sub>. The rate of change of each chemical family is given by Eqs. (5), (6), and (7): (Cariolle et al., 2009).

$$\frac{d([\text{O}] + [\text{O}_3])}{dt} = +k_2 \cdot [\text{NO}_2] - k_3 \cdot [\text{NO}] \cdot [\text{O}_3], \\ - k_5 \cdot [\text{O}] \cdot [\text{NO}_2] - 2 \cdot k_6 \cdot [\text{O}_3] \cdot [\text{O}], \quad (5)$$

$$\frac{d([\text{O}] + [\text{O}_3] + [\text{NO}_2])}{dt} = -2 \cdot k_5 \cdot [\text{O}] \cdot [\text{NO}_2], \\ - 2 \cdot k_6 \cdot [\text{O}_3] \cdot [\text{O}] \quad (6)$$

$$\frac{d([\text{NO}] + [\text{NO}_2])}{dt} = 0, \quad (7)$$

where  $k_i$  corresponds to the rate constants for the  $R_i$  reactions.

Thus, two processes occur to O<sub>3</sub> in the plume during the daytime. On short timescales O<sub>x</sub> is conserved. Lightning emissions of NO in the plume are converted into NO<sub>2</sub>, but as NO<sub>2</sub> is in O<sub>x</sub> family, there is net conservation of O<sub>x</sub>. However, on long timescales O<sub>x</sub> can be destroyed through the reaction of O with NO<sub>2</sub> and O<sub>3</sub>. Both of these processes need to be considered.

The first regime (regime I) occurs at low concentrations of NO<sub>x</sub> (relative to O<sub>3</sub>). Under these conditions the Reaction (R5) is slow. There is the rapid equilibrium between NO,

NO<sub>2</sub>, and O<sub>3</sub> (Reactions R2, R3, and R4). As a consequence, O<sub>3</sub> is converted into NO<sub>2</sub> and can be restored later after dilution of the plume depending on the balance between NO and NO<sub>2</sub> on a large scale (Cariolle et al., 2009). Overall, O<sub>x</sub> is conserved. In this regime emitted NO reacts with the available O<sub>3</sub> until the NO to NO<sub>2</sub> ratio in the plume reaches that in the background. Thus, the impact on the O<sub>3</sub> background concentration is to reduce it by the number of molecules of NO emitted multiplied by the background NO<sub>2</sub> to NO<sub>x</sub> ratio. The effect of the first regime on the ozone burden is expressed by Eq. (8).

$$\frac{\partial \overline{r_{\text{O}_3}}}{\partial t} + \langle F_{\text{O}_3} \rangle = -\frac{1}{\tau} \cdot \overline{r_{\text{LNO}_x}} \cdot \alpha_{\text{NO}_x} \cdot \left( \frac{\overline{\text{NO}_2}}{\overline{\text{NO}_x}} - E \right) \cdot \delta + L_{\text{ss}}, \quad (8)$$

where  $\overline{r_{\text{O}_3}}$  is the mixing ratio of O<sub>3</sub> (in ppb) in the model grid,  $E$  is the  $\frac{\text{NO}_2}{\text{NO}_x}$  ratio in the initial emissions,  $\delta$  is equal to 1 during the day and 0 during the nighttime, and  $L_{\text{ss}}$  is the sources and sinks of ozone, such as photochemical production, transport from the stratosphere, surface deposition, photolysis reactions, and photochemical destruction.

The second regime (regime II) occurs at high concentrations of NO<sub>x</sub> (relative to O<sub>3</sub>). Under these conditions the rate of R5 is large. The nonlinear chemical interactions between NO<sub>x</sub> and O<sub>3</sub> occur with different rates than in the background atmosphere. To account for this, Cariolle et al. (2009) introduced an effective reaction rate constant ( $K_{\text{eff}}$ ), which is related to the production or the destruction of the odd oxygen (O<sub>x</sub>) within the plume.  $K_{\text{eff}}$  is expressed by Eq. (9).

$$K_{\text{eff}} = \frac{\int_{t_0}^{T_1} \left( \int_{V_p} K \cdot r_{\text{NO}_x}^p \cdot r_{\text{O}_3}^p \cdot dV_p \right) \cdot dt}{\overline{r_{\text{O}_3}} \cdot \int_{t_0}^{T_1} \left( \int_{V_p} r_{\text{NO}_x}^p \cdot dV_p \right) \cdot dt}, \quad (9)$$

where  $r_{\text{NO}_x}^p$  and  $r_{\text{O}_3}^p$  are the mixing ratios of nitrogen oxides and ozone within the plume,  $\overline{r_{\text{O}_3}}$  is the background ozone mixing ratio averaged in the model grid, and  $K$  is the rate of NO<sub>x</sub>–O<sub>3</sub> reaction within the plume.

The analysis of the chemical reactions related to the two regimes shows that  $[\text{O}_3] \gg [\text{O}]$  and  $k_5 \times [\text{NO}_2]$  is more efficient than  $k_6 \times [\text{O}_3]$  as a sink for O<sub>x</sub> (Cariolle et al., 2009). Thus, Eq. (6) is simplified to give Eq. (10).

$$\frac{d([\text{O}_3] + [\text{NO}_2])}{dt} = -2 \cdot k_5 \cdot [\text{O}] \cdot [\text{NO}_2] \quad (10)$$

Consequently,  $K_{\text{eff}}$  can be simplified to Eq. (11).

$$K_{\text{eff}} = \frac{2 \cdot \left( \int^T k_5 \cdot \text{O} \cdot \text{NO}_2 \cdot dt \right)}{(\text{NO}_x \cdot \int^T \text{O}_x \cdot dt)} \quad (11)$$

The calculation of  $K_{\text{eff}}$  is detailed in Sect. 3.2.4. Considering the two regimes related to the sub-grid plume chemistry, the ozone burden is described by Eq. (12) during the daytime and nighttime. Note that during the nighttime there is no direct

impact on its burden due to the ozone plume chemistry as  $\delta = 0$ . Only indirect effects are expected from NO<sub>y</sub> chemistry.

$$\frac{\partial \overline{r_{O_3}}}{\partial t} + \langle F_{O_3} \rangle = -\frac{1}{\tau} \cdot \overline{r_{LNO_x}} \cdot \alpha_{NO_x} \cdot \left( \frac{\overline{NO_2}}{\overline{NO_x}} - E \right) \cdot \delta - K_{eff} \cdot \overline{r_{LNO_x}} \cdot \rho \cdot \alpha_{NO_x} \cdot \overline{r_{O_3}} \cdot \delta + L_{ss} \quad (12)$$

In addition, we consider the conversion of NO<sub>x</sub> into HNO<sub>3</sub> within the plume. This conversion takes place in two different ways depending on the day or night atmospheric conditions. During the day, NO<sub>2</sub> reacts primarily with OH to give HNO<sub>3</sub> directly, and it is characterized by the coefficient  $\beta_1$ , while during the nighttime the conversion of NO<sub>x</sub> to HNO<sub>3</sub> occurs mainly through the N<sub>2</sub>O<sub>5</sub> formation followed by a heterogeneous hydrolysis reaction, which corresponds to  $\beta_2$ . In other words, the  $\beta$  coefficients are the molar fractions of NO<sub>x</sub> converted to HNO<sub>3</sub> within the plume. These two fractions are unitless.

In summary, the equation system solved on a large scale by the CTM for a lightning NO<sub>x</sub> source is detailed by Eqs. (13), (14), and (15).

$$\frac{\partial \overline{r_{NO_x}}}{\partial t} + \langle F_{NO_x} \rangle = +\frac{1}{\tau} \cdot \overline{r_{LNO_x}} \cdot (1 - \beta_1 \cdot \delta - \beta_2 \cdot (1 - \delta)) \cdot \alpha_{NO_x} + L_{ss}, \quad (13)$$

$$\frac{\partial \overline{r_{HNO_3}}}{\partial t} + \langle F_{HNO_3} \rangle = +\frac{1}{\tau} \cdot \overline{r_{LNO_x}} \cdot (\beta_1 \cdot \delta + \beta_2 \cdot (1 - \delta)) \cdot \alpha_{NO_x} + L_{ss}, \quad (14)$$

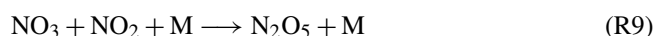
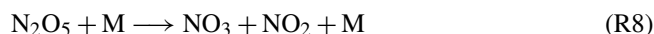
$$\frac{\partial \overline{r_{O_3}}}{\partial t} + \langle F_{O_3} \rangle = -\left( \frac{1}{\tau} \cdot \left( \frac{\overline{NO_2}}{\overline{NO_x}} - E \right) + K_{eff} \cdot \overline{r_{O_3}} \cdot \rho \right) \cdot \overline{r_{LNO_x}} \cdot \alpha_{NO_x} \cdot \delta + L_{ss}, \quad (15)$$

where  $\overline{r_{NO_x}}$ ,  $\overline{r_{HNO_3}}$ , and  $\overline{r_{O_3}}$  correspond to the NO<sub>x</sub>, HNO<sub>3</sub>, and O<sub>3</sub> mixing ratios averaged over the grid cell of the model, respectively.

In this study, the tropospheric chemistry and especially the LNO<sub>x</sub> plume chemistry is considered both during the daytime and nighttime since no reactions are initiated during the day. The chemical interactions during the night correspond mainly to the reactions of O<sub>3</sub> and O with NO and NO<sub>2</sub> as well as the NO<sub>x</sub> deactivation and the chemistry of the nitrogen reservoir species (here, HNO<sub>3</sub> and N<sub>2</sub>O<sub>5</sub>) and the nitrate radical (NO<sub>3</sub>). NO<sub>3</sub> is the main oxidant in night conditions and is produced from the slow oxidation of NO<sub>2</sub> by O<sub>3</sub> (Reaction R7).

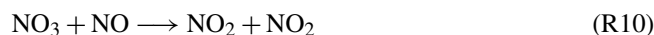


The other dominant source of NO<sub>3</sub> is the destruction of N<sub>2</sub>O<sub>5</sub> (Reaction R8), but as N<sub>2</sub>O<sub>5</sub> is formed from NO<sub>3</sub> (Reaction R9), the two species act in a coupled manner.



As mentioned previously, N<sub>2</sub>O<sub>5</sub> is a determining species for the tropospheric chemistry during the nighttime, allow-

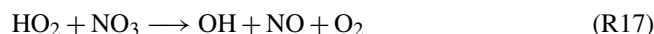
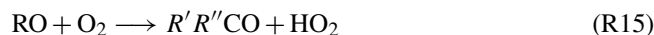
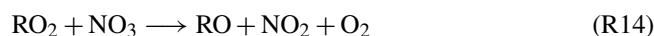
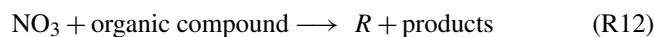
ing the HNO<sub>3</sub> formation by the heterogeneous reaction on the particle surface (aerosols and ice crystals). During the day, NO<sub>3</sub> rapidly undergoes photolysis to produce NO or NO<sub>2</sub>. In addition, NO<sub>3</sub> reacts very quickly with NO, which is more concentrated during the daytime than during the nighttime (Reaction R10) but NO<sub>3</sub> is very low during the daytime. However, this reaction can take place during the night, especially for a plume characterized by high NO mixing ratios (like a plume from lightning emissions) which is transported both during the day and at night.



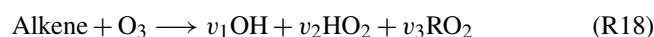
Furthermore, the nitrate radical can potentially react with VOCs. The reaction of the unsaturated hydrocarbons such as isoprene, butenes, and monoterpenes with NO<sub>3</sub> leads to the HNO<sub>3</sub> formation (Monks, 2005) (Reaction R11).



Considering NO<sub>3</sub> reaction with alkenes, an additional mechanism is found initiating a complex chemistry allowing the formation of NO<sub>2</sub> or organic nitrates (Monks, 2005). Finally, NO<sub>3</sub> can initiate VOC oxidation via peroxy radical production (Reaction R12). That way, it can be involved as a chain propagator (Reactions R13 to R17).



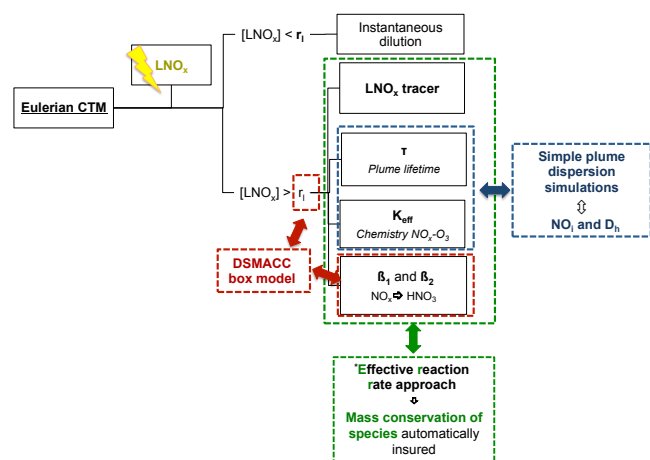
The reactions of HO<sub>2</sub> with ozone (R16) or NO<sub>3</sub> (Reaction R17) imply OH production. Also, the reaction of ozone with alkenes allows the formation of OH during the night (Reaction R18) (Aumont, 2005).



Reaction (R18) occurs when ozone concentrations remain sufficiently high in night conditions, in other words for polluted atmosphere.

In this context, we consider different values during the daytime and nighttime for the plume lifetime, the effective reaction rate constant, and the fraction of NO<sub>x</sub> conversion into HNO<sub>3</sub> within the plume. Distinguishing day and night chemistry is linked with the fluctuation of the critical  $r_1$  value (below which the sub-grid plume chemistry is negligible), depending on atmospheric conditions. Therefore, if  $r_1$  changes with sunlight, the plume lifetime also changes. Note that except for the  $\beta_2$  fraction, this night chemistry is not considered by the initial plume approach developed by Carille et al. (2009), which considers NO<sub>x</sub> plumes from aircraft exhausts only during the daytime.





**Figure 1.** Diagram of the lightning NO<sub>x</sub> plume parameterization based on the effective reaction rate approach. The arrows link the parameters to their estimate approach. The red boxes are for the parameters estimated with the DSMACC model and the blue boxes are related to the parameters calculated with the simple plume dispersion model. Finally, the green boxes show the effective reaction rate approach in the GEOS-Chem CTM.

Figure 1 summarizes all elements which define the plume approach and how it has been adapted and implemented into the model.

### 3.2 Parameter calculations for lightning NO<sub>x</sub> emissions

In order to reproduce more accurately the lightning NO<sub>x</sub> sub-grid chemistry, some points should be considered: (i) the latitude (NO<sub>x</sub> emissions by lightning are higher in the tropics than in the midlatitudes); (ii) the sunlight conditions (day and night), which impact photochemistry and heterogeneous chemistry; (iii) the plume evolution with its own physical characteristics (the lifetime and dispersion properties); and finally (iv) chemical interactions within the plume related to highly elevated NO<sub>x</sub> concentrations relative to the background. In the following section, physical and chemical characteristics of the plume associated with a lightning NO<sub>x</sub> source have been defined.

#### 3.2.1 Dynamical conditions

The horizontal diffusion coefficient ( $D_h$ ) is a key parameter of the atmospheric dynamical conditions in determining the dispersion of the lightning NO<sub>x</sub> plume.  $D_h$  is used as the dispersion constraint for the simple plume dispersion simulations carried out in order to estimate the plume lifetime and the effective reaction rate constant. The diffusion coefficient was determined in two different ways. A first estimate of the horizontal diffusion was performed by running the 3-D mesoscale Meso-NH model. Then, the  $D_h$  coefficient was calculated using in situ measurements in a thunderstorm anvil.

The Meso-NH mesoscale model was used (see Sect. 2.2) to investigate  $D_h$ . A simple convective cell forced by a warm bubble and initialized by a radiosounding at the simulation start was run as an ideal case. Simulations were realized for a domain of 24 km in the two horizontal directions, and the grid horizontal resolution is  $\Delta x = \Delta y = 1$  km and  $\Delta z = 500$  m. The convective cell is located at 43.29° N latitude and 0° longitude (Klemp and Wilhelmson, 1978). Simulations of 6 h were performed, allowing the complete development and the dissipation of the convective cell.  $D_h$  was calculated within the anvil using the mixing length diagnostic variable, hereafter denoted  $L$ , as described by Eq. (16) (Cuxart et al., 1999).

$$D_h = \frac{2}{3} \times \frac{L}{4} \times \exp\left(\frac{1}{2}\right) \quad (16)$$

At the mature stage of the cell,  $D_h$  was calculated as  $100 \text{ m}^2 \text{ s}^{-1}$  within the upper levels of the convective cell (i.e., in the anvil, defined empirically).

In addition to the modeling estimate, we used in situ measurements to calculate  $D_h$ . Turbulence measurements were performed by a B-757 commercial aircraft along a flight from the west of Kansas to the north of Missouri and corresponding to a trajectory of more than 500 km (Trier and Sharman, 2008). These in situ measurements were carried out from 07:00 to 10:00 UTC on 17 June 2005, during the development of a mesoscale convective system (MCS). This MCS is associated with a turbulence event characterized by the measurement of the atmospheric eddy dissipation rate ( $\epsilon$ ) and the turbulence kinetic energy (TKE) above and within the cloud anvil. The higher values of  $\epsilon$  ( $\epsilon^{1/3} \sim 0.4 \text{ m}^{2/3} \text{ s}^{-1}$ ) were recorded between 11.3 and 11.6 km altitude, corresponding to the cloud anvil levels. In addition, for this MCS, the TKE was about  $1 \text{ m}^2 \text{ s}^{-2}$  at the locations of the highest  $\epsilon$  values.

According to these observations, the turbulent diffusivity (Eq. 17) was estimated above the anvil of the MCS ([http://www.ral.ucar.edu/projects/turb\\_char/](http://www.ral.ucar.edu/projects/turb_char/)), such that  $D_h > 0.1 \text{ m}^2 \text{ s}^{-2}$ . Then,  $D_h$  was calculated within the anvil such that  $D_h = 15 \text{ m}^2 \text{ s}^{-1}$  using the same formulation (Eq. 17). This last estimate seems to be the most common value compared to the diffusion coefficient value of  $20 \text{ m}^2 \text{ s}^{-1}$  used by Cariolle et al. (2009), close to the tropopause level and the  $D_h$  value calculated for contrails ( $15 \text{ m}^2 \text{ s}^{-1}$ ) in the upper troposphere (Knollenberg, 1972).

$$D_h = \frac{(\text{TKE})^2}{\epsilon} \quad (17)$$

The  $D_h$  estimate using the Meso-NH model is high compared to the results from measurements, corresponds to the upper limit of the calculated diffusion coefficients, and may be associated with the turbulence in the convective cloud. However, it is important to note that usually most numerical simulations are performed with 1-D turbulence models. What is interesting in the use of Meso-NH in this study is that the 3-D



turbulence is solved. This simulation provides an additional estimate of  $D_h$ , allowing a comparison with the calculation from in situ measurements. Moreover, studies on the diffusivity in cloud anvils are uncommon. It is necessary to conduct additional work in the future on that issue, again constrained with new in situ measurements of the atmospheric turbulence in the anvil.

It is important to note that the 3-D turbulence is not solved online in the GEOS-Chem model because of the fine scale characterizing this process but is prescribed by the GEOS-5 met fields. Therefore, the global variability of  $D_h$  is not calculated by the CTM, and it is beyond the scope of this study.

In order to cover all horizontal diffusivity estimates discussed in this section, values of 0.1, 15, and 100 m<sup>2</sup> s<sup>-1</sup> were used. The horizontal coefficient is constant for all lightning NO<sub>x</sub> plumes which are considered in the GEOS-Chem model. Hereafter, the results are detailed for the central value  $D_h = 15 \text{ m}^2 \text{ s}^{-1}$ . Sensitivity tests depending on the uncertainty associated with the parameter estimate are performed and presented later in Sect. 4.3.

### 3.2.2 The NO<sub>x</sub> critical plume content ( $r_l$ )

The  $r_l$  critical value is the NO<sub>x</sub> mixing ratio within the undiluted phase of the plume below which the nonlinear chemistry can be neglected (Sect. 3.1). It has been estimated using the 0-D DSMACC box model (Sect. 2.2). Initial conditions for simulations carried out with the DSMACC box model are from outputs of the GEOS-Chem model. In particular, initial atmospheric parameters and atmospheric background concentrations of species correspond to the average of the GEOS-Chem outputs (i) from 8 to 11 km, (ii) for two latitude regions (tropics and midlatitudes), and (iii) for the year 2006 (Table 1). The altitude range refers to the detrainment region estimated by GEOS-Chem using the GEOS-5 met fields (Sect. 2.1) both in the tropics and in the midlatitudes. Note that this range can vary depending on the met fields and the convection parameterization. In addition, the LNO<sub>x</sub> plume parameterization should have an impact outside of this altitude range, mainly between 6 and 12 km, but to a lesser extent.

In order to focus on chemistry interactions only between chemical species of interest and on removing the mixing influence and sunlight fluctuations, short simulations (i.e., 1 h each) were run with the DSMACC model. The effects of the day or night conditions were carefully considered, carrying out separate simulations during the daytime and nighttime. Simulations were run for a large range of initial NO mixing ratios from 0.01 ppb to 1 ppm. The  $r_l$  value is defined from the NO value, for which the  $\frac{\partial \text{O}_x}{\partial t}$  trend is perturbed. In other words,  $r_l$  is associated with the second derivative of O<sub>x</sub>, i.e., the curve optimums on Fig. 2. The  $r_l$  threshold was defined to be 0.1 and 0.25 ppb during the day and at night for mid-

latitudes and 0.1 and 0.75 ppb during the day and at night in the tropics (Fig. 2).

Note that the midlatitudes and the tropics were separated because of the large differences in LNO<sub>x</sub> emissions between the two regions in terms of the number of flashes in a particular convective cell, which is higher in the tropics according to the LIS-OTD climatologies (Christian et al., 2003). This last point is important for the plume lifetime estimate detailed in the following section.

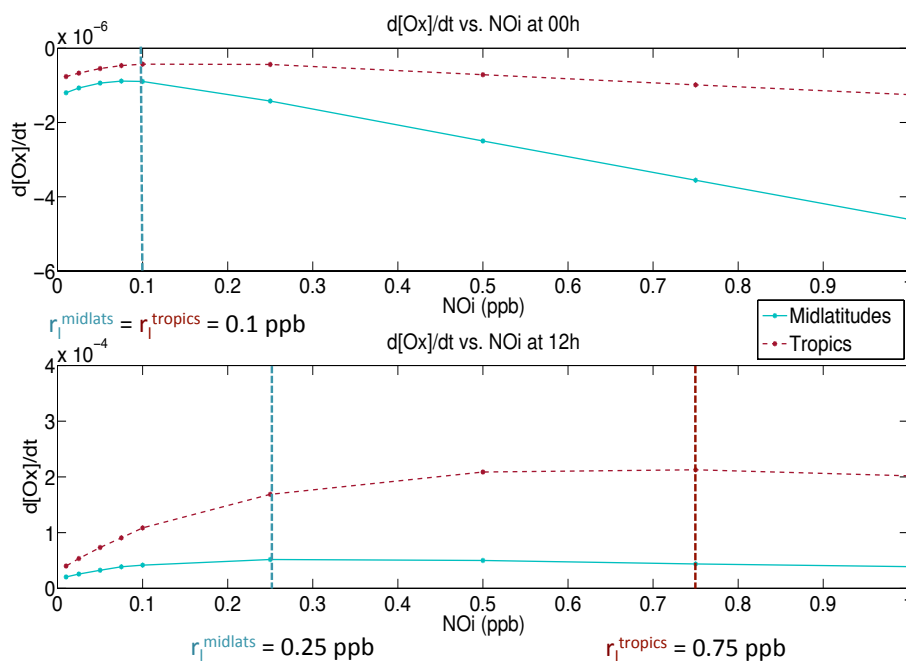
### 3.2.3 The plume lifetime $\tau$

The plume lifetime ( $\tau$ ) depends directly on (i) the initial NO pulse from lightning emissions, (ii) the  $r_l$  critical value, and (iii) the diffusion properties of the atmosphere. The plume lifetime also depends on the initial size of the plume. Here we use a width of 500 m in order to include an ensemble of spikes on the cloud scale (i.e., each plume is defined from several electrical discharges on a convective cell scale).  $\tau$  is crucial for the physical description of the NO<sub>x</sub> plumes, and it has been computed by carrying out dispersion simulations of a simple plume assumed to be cylindrical. In this model, the standard atmospheric conditions are represented by temperature, pressure, and species concentrations of the background atmosphere, which are similar to the initial conditions used for the DSMACC simulations. As a reminder, initial conditions are from GEOS-Chem outputs averaged (i) from 8 to 11 km, (ii) for two latitude regions (tropics and midlatitudes), and (iii) for the year 2006 (Table 1). Simulations are initialized by an NO pulse from lightning emissions (hereafter denoted NO<sub>i</sub>), and the plume dispersion depends on the  $D_h$  value estimated in Sect. 3.2.1.

The initial tracer concentration NO<sub>i</sub> related to lightning NO emissions on the scale of a convective cell (gathering several flashes together) in the midlatitudes was defined according to previous aircraft measurement campaigns. In particular, the STERAO (Stratospheric–Tropospheric experiment: radiation, aerosols, and ozone) campaign recorded NO spikes of a magnitude from 1 to 10 ppb related to lightning activity in thunderstorms and occurring from 9 to 10 July 1996 over northern Colorado (Dye et al., 2000; Stith et al., 1999). Lange et al. (2001) measured NO spikes of 3.5 ppb during the STREAM (Stratosphere–Troposphere Experiment by Aircraft Measurements) campaign associated with a matured storm over Ontario. Several peaks of NO mixing ratios from 0.7 to 6 ppb were also observed during EULINOX (European Lightning Nitrogen Oxides Project; Huntrieser et al., 2002) over Germany in July 1998. The LINOX (Lightning NO<sub>x</sub>) aircraft campaign recorded NO spikes from 0.75 to 1.25 ppb (Huntrieser et al., 1998) related to thunderstorms over Europe on 30 July 1996. From these studies, the NO concentration associated with the electrical activity in thunderstorms occurring over the midlatitudes was determined as  $\text{NO}_i^{\text{mean, Midlats}} = 3.4 \text{ ppb}$  ( $\text{NO}_i^{\text{min, Midlats}} = 0.7 \text{ ppb}$  and  $\text{NO}_i^{\text{max, Midlats}} = 10 \text{ ppb}$ ). Because there are much fewer

**Table 1.** The initial atmospheric parameters and background concentrations of chemical species from GEOS-Chem outputs for the DSMACC box model simulations.

	Temp.	Press.	O <sub>3</sub>	NO	NO <sub>2</sub>	HNO <sub>3</sub>	HNO <sub>4</sub>	PAN	N <sub>2</sub> O <sub>5</sub>	CO
Units	(K)	(hPa)	(ppb)	(ppb)	(ppb)	(ppb)	(ppb)	(ppb)	(ppt)	(ppb)
Midlatitudes	228	313	67	0.04	0.01	0.15	0.02	0.1	2	94
Tropics	240	313	26	0.03	0.003	0.02	0.006	0.03	2.3	93
	OH	HO <sub>2</sub>	H <sub>2</sub> O <sub>2</sub>	CH <sub>2</sub> O	CH <sub>4</sub> O <sub>2</sub>	C <sub>3</sub> H <sub>8</sub>	C <sub>5</sub> H <sub>8</sub>	C <sub>2</sub> H <sub>4</sub> O	C <sub>3</sub> H <sub>6</sub> O	
Units	(ppb)	(ppt)	(ppt)	(ppb)	(ppb)	(ppb)	(ppb)	(ppb)	(ppb)	
Midlatitudes	0.2	4	0.4	0.06	0.1	0.47	0	7.5	4	
Tropics	0.06	6	0.34	0.03	0.17	0.13	7.5	7.5	4	

**Figure 2.**  $r_l$  critical value and odd oxygen trends from DSMACC box model simulations for midlatitudes (solid line) and tropics (dotted line) at midnight (upper panel) and at midday (bottom panel).

LNO<sub>x</sub> measurements in the tropics and in order to be consistent with the LNO<sub>x</sub> emissions defined in the GEOS-Chem model, the ratio  $R_{\text{LNO}_x} = \frac{\text{LNO}_x^{\text{Midlatitudes}}}{\text{LNO}_x^{\text{Tropics}}}$  was defined as in the CTM. During the year 2006, the relative midlatitude and tropics LNO<sub>x</sub> contribution was about  $R_{\text{LNO}_x} = 0.33$ . This result is in agreement with higher LNO<sub>x</sub> emissions in these regions rather than in the midlatitudes. The value of the NO mixing ratio injected by lightning into the tropics was estimated as  $\text{NO}_i^{\text{mean, Tropics}} = 10.2$  ppb ( $\text{NO}_i^{\text{min, Tropics}} = 2.8$  ppb and  $\text{NO}_i^{\text{max, Tropics}} = 29.7$  ppb).

Once the NO<sub>i</sub> estimate was completed, the calculation of the plume lifetime was achieved using the detailed formulation given in Sect. 3.1.1. The results for  $\tau$  are summarized in Table 2. Hereafter, the results are detailed for  $\text{NO}_i^{\text{mean}}$  in Sect. 4, and sensitivity tests are carried out using all NO<sub>i</sub>

values for the midlatitudes and the tropics (Sect. 5). Model calculations for  $\text{NO}_i^{\text{mean}}$  and  $D_h = 15 \text{ m}^2 \text{ s}^{-1}$  provide a minimum plume lifetime of 3 (6) h for the midlatitudes and maximum plume lifetime of 9 (21.3) h for the tropics during the daytime (nighttime).

### 3.2.4 The effective reaction rate constant ( $K_{\text{eff}}$ )

The nonlinear chemistry within the plume has been considered in calculating the effective reaction rate constant ( $K_{\text{eff}}$ ), which is used to compute the formation of the secondary species (O<sub>x</sub> and HNO<sub>3</sub>) within the plume.  $K_{\text{eff}}$  is associated with the evolution of odd oxygen depending on the O and O<sub>3</sub> reactions with NO<sub>2</sub> and NO, and also on the NO<sub>x</sub> activation (day) or deactivation (night) with the HNO<sub>3</sub>, N<sub>2</sub>O<sub>5</sub>, and PAN chemistry. Note that in the case of lightning emissions,

**Table 2.** The plume lifetime  $\tau$  (hours) calculated for midlatitudes and tropics depending on the initial NO mixing ratio injected by lightning emissions ( $\text{NO}_i$ , ppb) and the horizontal diffusion coefficient ( $D_h$ ,  $\text{m}^2 \text{s}^{-1}$ ) for day (upper part of table) and night conditions (bottom part of table).

$\tau$ (hours)	Day			Tropics		
	Midlatitudes					
$\text{NO}_i$ (ppb)	0.7	3.4	10	2.8	10	29.7
$D_h = 0.1$ ( $\text{m}^2 \text{s}^{-1}$ )	1.55	8.14	23.9	4.40	23.1	67.9
$D_h = 15$ ( $\text{m}^2 \text{s}^{-1}$ )	0.1	3.17	18.6	0.27	8.90	52.8
$D_h = 100$ ( $\text{m}^2 \text{s}^{-1}$ )	0.01	0.47	4.17	0.04	1.32	11.7

$\tau$ (hours)	Night			Tropics		
	Midlatitudes					
$\text{NO}_i$ (ppb)	0.7	3.4	10	2.8	10	29.7
$D_h = 0.1$ ( $\text{m}^2 \text{s}^{-1}$ )	1.62	8.19	24.1	4.74	23.4	68.5
$D_h = 15$ ( $\text{m}^2 \text{s}^{-1}$ )	0.31	6.19	22	2.77	21.3	66.4
$D_h = 100$ ( $\text{m}^2 \text{s}^{-1}$ )	0.05	1.23	10.6	0.43	10.5	55.4

other species like VOCs, HO<sub>x</sub>, and H<sub>2</sub>O may be uplifted into the convective region. However, we assumed that the OPE is mainly controlled by NO<sub>x</sub> in the upper troposphere, as previously showed by Sauvage et al. (2007b). Therefore, the  $K_{\text{eff}}$  calculation is here mainly dependent on NO<sub>x</sub> content. Future studies should try to investigate this issue for lightning emissions mixed with strong surface emissions in order to sharpen our parameterization.

$K_{\text{eff}}$  is calculated according to Eq. (11) of Sect. 3.1.2 using the same simple plume dispersion simulations as those carried out to define the plume lifetime (Sect. 3.2.3).

Results for  $K_{\text{eff}}$  are summarized in Table 3. Model calculations using  $\text{NO}_i^{\text{mean}}$  and  $D_h = 15 \text{ m}^2 \text{s}^{-1}$  give a  $K_{\text{eff}}$  value of  $5.49 \times 10^{-19} \text{ molecules}^{-1} \text{s}^{-1} \text{cm}^3$  ( $4.55 \times 10^{-19} \text{ molecules}^{-1} \text{s}^{-1} \text{cm}^3$ ) in the midlatitudes and  $3.64 \times 10^{-19} \text{ molecules}^{-1} \text{s}^{-1} \text{cm}^3$  ( $2.98 \times 10^{-19} \text{ molecules}^{-1} \text{s}^{-1} \text{cm}^3$ ) in the tropics during the daytime (during the nighttime).

Our  $K_{\text{eff}}$  estimates are smaller than those calculated by Cariolle et al. (2009) for the plume chemistry related to aircraft exhausts. In this previous work,  $K_{\text{eff}}$  varies from 1.0 to  $4.2 \times 10^{-18} \text{ molecules}^{-1} \text{s}^{-1} \text{cm}^3$ , with a mean value close to  $3 \times 10^{-18} \text{ molecules}^{-1} \text{s}^{-1} \text{cm}^3$  depending on the NO<sub>x</sub> loading. The very low value for  $K_{\text{eff}}$  indicates that the plume parameterization implies a delay in the production of ozone on a large scale rather than its destruction within the plume.

### 3.2.5 The fractions of NO<sub>x</sub> conversion to HNO<sub>3</sub> ( $\beta_1$ and $\beta_2$ )

The fractions  $\beta_1$  and  $\beta_2$  represent the NO<sub>x</sub> conversion to HNO<sub>3</sub> within the plume during the daytime and nighttime, respectively. They were computed using the DSMACC box model.

The  $\beta_1$  coefficient was calculated for day conditions depending mainly on the OH concentration. The conversion of NO<sub>x</sub> into HNO<sub>3</sub> during the nighttime ( $\beta_2$  coefficient) is related to the heterogeneous reaction of N<sub>2</sub>O<sub>5</sub> and so depends on particle (aerosols and ice crystals) concentration and particles' lifetime. This is directly linked with the surface density and the radius of particles in the anvil region of thunderstorms, which is highly uncertain. We defined these values using in situ measurements. The surface area ( $S_T$ ) and the radius ( $R$ ) for aerosols are defined such that  $S_T = 0.28 \text{ m}^{-1}$  and  $R = 1 \mu\text{m}^{-1}$  (Huntrieser et al., 2002); for ice, these are  $S_T = 0.03 \text{ m}^{-1}$  and  $R = 30 \mu\text{m}^{-1}$  (Knollenberg et al., 1993). In addition, the reaction probabilities of NO<sub>x</sub> on aerosols and ice crystals  $\gamma_{\text{N}_2\text{O}_5}^{\text{aerosols}} = 0.02$  (Evans and Jacob, 2005) and  $\gamma_{\text{N}_2\text{O}_5}^{\text{ice}} = 0.03$  (Sander et al., 2006), respectively, were used for our box model simulations. These values correspond to the probability that an N<sub>2</sub>O<sub>5</sub> molecule impacting an aerosol or an ice crystal surface reacted. The results for  $\beta_1$  and  $\beta_2$  coefficients are summarized in Table 4.

The estimate of the  $\beta_1$  fraction does not show significant variation, neither between latitudes regions nor depending on NO<sub>i</sub>. The minimum  $\beta_1$  value is  $1.34 \times 10^{-4}$  for the tropical regions and  $\text{NO}_i^{\text{min}}$ , and the maximum  $\beta_1$  value is  $1.88 \times 10^{-4}$  for the midlatitudes and  $\text{NO}_i^{\text{max}}$ . The study of production and destruction rates for day conditions taking into account all reactions pathways (not shown here) demonstrates that the production of HNO<sub>3</sub> during the day is mainly determined by the reaction of NO<sub>3</sub> with formaldehyde (HCHO) and acetaldehyde (CH<sub>3</sub>CHO). Surprisingly, the HNO<sub>3</sub> formation via the NO<sub>2</sub>+OH reaction seems to be less efficient. This result may be explained by the low initial concentrations of OH used for the DSMACC simulations, and it is in agreement with the small  $\beta_1$  values. The averaged  $\beta_2$  coefficient is higher by a factor 10 compared to  $\beta_1$  with a minimum value of  $0.24 \times 10^{-3}$  in the tropics for  $\text{NO}_i^{\text{max}}$  and a maximum estimate of  $14.4 \times 10^{-3}$  in the midlatitudes for  $\text{NO}_i^{\text{min}}$ . The analysis of the production and the destruction rates for night conditions taking into account all reaction pathways shows that the predominant reaction in the HNO<sub>3</sub> evolution is N<sub>2</sub>O<sub>5</sub>+H<sub>2</sub>O (or the heterogeneous reaction on the aerosol and ice crystal surfaces).

## 4 Results: CTM simulations

In this section, the effects of the lightning NO<sub>x</sub> plume parameterization, i.e., the influence of the sub-grid processes related to lightning emissions, on the NO<sub>x</sub> and O<sub>3</sub> tropospheric distributions on a large scale are evaluated. Then, the parameterization sensitivity to the initial NO mixing ratio injected by lightning ( $\text{NO}_i$ ) and the  $D_h$ ,  $\beta_1$ , and  $\beta_2$  coefficients are analyzed to quantify the variability of the results regarding the plume-in-grid parameter calculations.

**Table 3.** The effective reaction rate constant  $K_{\text{eff}}$  ( $10^{-19}$  molecules<sup>-1</sup> s<sup>-1</sup> cm<sup>3</sup>) in the midlatitudes and tropics depending on the initial NO mixing ratio injected by lightning emissions ( $\text{NO}_i$ , ppb) and the horizontal diffusion coefficient ( $D_h$ , m<sup>2</sup> s<sup>-1</sup>) for day (upper part of table) and night conditions (bottom part of table).

$K_{\text{eff}}$ ( $10^{-19}$ molecules <sup>-1</sup> s <sup>-1</sup> cm <sup>3</sup> )	Day					
	Midlatitudes			Tropics		
$\text{NO}_i$ (ppb)	0.7	3.4	10	2.8	10	29.7
$D_h = 0.1$ (m <sup>2</sup> s <sup>-1</sup> )	1.28	1.24	1.51	0.77	1.2	1.83
$D_h = 15$ (m <sup>2</sup> s <sup>-1</sup> )	8.44	5.49	5.43	7.79	3.64	4.13
$D_h = 100$ (m <sup>2</sup> s <sup>-1</sup> )	12.1	16.4	14.4	23	19.8	13

$K_{\text{eff}}$ ( $10^{-19}$ molecules <sup>-1</sup> s <sup>-1</sup> cm <sup>3</sup> )	Night					
	Midlatitudes			Tropics		
$\text{NO}_i$ (ppb)	0.7	3.4	10	2.8	10	29.7
$D_h = 0.1$ (m <sup>2</sup> s <sup>-1</sup> )	1.28	1.24	1.51	0.77	1.10	1.83
$D_h = 15$ (m <sup>2</sup> s <sup>-1</sup> )	4.84	4.55	5.43	2.3	2.98	4.13
$D_h = 100$ (m <sup>2</sup> s <sup>-1</sup> )	7.36	8.39	6.73	6.45	3.94	5.16

**Table 4.** The fractions of NO<sub>x</sub> conversion into HNO<sub>3</sub> within the plume ( $\beta_1$  and  $\beta_2$ ) in the midlatitudes and tropics depending on the initial NO mixing ratio injected by lightning emissions ( $\text{NO}_i$ , ppb) and on particles for day (upper part of table) and night conditions (bottom part of table).

$\beta_1$ ( $10^{-4}$ )	Day					
	Midlatitudes			Tropics		
$\text{NO}_i$ (ppb)	0.7	3.4	10	2.8	10	29.7
Aerosols	2.53	3.34	3.45	2.51	2.95	2.6
Ice	0.23	0.3	0.3	0.2	0.23	0.3
Mean	1.38	1.8	1.88	1.34	1.59	1.47

$\beta_2$ ( $10^{-3}$ )	Night					
	Midlatitudes			Tropics		
$\text{NO}_i$ (ppb)	0.7	3.4	10	2.8	10	29.7
Aerosols	14.3	9.89	8	4.9	1.69	0.24
Ice	14.4	9.96	8.06	4.89	1.70	0.24
Mean	14.4	9.92	8.03	4.88	1.7	0.24

#### 4.1 Implementation of the LNO<sub>x</sub> plume parameterization

The implementation of the lightning NO<sub>x</sub> plume parameterization into the GEOS-Chem model requires specifying the system of continuity equations related to the plume chemistry solved on a large scale by the model (Sect. 3.1.2, Eqs. 13, 14, and 15). Lightning NO<sub>x</sub> emissions calculated in each grid box (in molecules cm<sup>-2</sup> s<sup>-1</sup>) by the model are directly used to compute the injection rate  $I$  (s<sup>-1</sup>) of NO at each chemical time step of the simulation. Then, we consider that  $\alpha_{\text{NO}_x} = 1$  in order to represent the mixing ratio of the undiluted fraction of NO<sub>x</sub> by the tracer ( $r_{\text{LNO}_x}$ ). Furthermore, lightning produces negligible quantities of NO<sub>2</sub> relative to NO, and therefore  $E$  is 0 in Eq. (15). Finally, the ratio  $\overline{\text{NO}_2/\text{NO}_x}$  is the

relative balance between NO and NO<sub>2</sub> in the diluted phase on a large scale reproduced by the model.

#### 4.2 Impact of LNO<sub>x</sub> emissions on the NO<sub>x</sub> and O<sub>3</sub> distributions

We perform a spin-up of 6 months (from July 2005 to January 2006) in order to obtain a steady state in the model after activation of the plume parameterization. Then, simulations were run for the entire year 2006. The transport and the convection time steps are 15 min and the emissions and the chemical time steps are 30 min.

In the following, “standard simulation” refers to a simulation with standard lightning NO<sub>x</sub> emissions, i.e., those instantaneously diluted in a grid cell, while “modified simulation” refers to a simulation considering the plume parameterization and then the sub-grid chemistry. Note that the modified simulation was run using mean values for the initial NO mixing ratio ( $\text{NO}_i^{\text{mean, Midlats}} = 3.4$  ppb and  $\text{NO}_i^{\text{mean, Tropics}} = 10.2$  ppb) and  $D_h = 15$  m<sup>2</sup> s<sup>-1</sup>. The base case (BC) experiment corresponds to the standard simulation minus the standard simulation without lightning NO<sub>x</sub> emissions. The P1 experiment corresponds to the modified simulation minus the standard simulation without lightning NO<sub>x</sub> emissions. The P2 experiment is the same as the P1 experiment but without considering the nitrification mechanism in the modified simulation (i.e.,  $\beta_1 = \beta_2 = 0$ ). In addition, sensitivity tests were performed for P1 defined by the modified simulation using the minimum and the maximum values for  $D_h$ ,  $\text{NO}_i$ , and the  $\beta_1$  and  $\beta_2$  coefficients. All experiments are summarized in Table 5.

Lightning emissions rates and the associated LNO<sub>x</sub> tracer distributions are first discussed; then the effects of the implementation of the plume parameterization (P1) compared to

**Table 5.** Values of the parameters for the plume parameterization corresponding to the experiments P1 and P2.

Parameters		Experiments											
		P1						P2					
$D_h$ (m <sup>2</sup> s <sup>-1</sup> )		0.1			15			100			15		
NO <sub>i</sub> (ppb)	Min	Mean	Max	Min	Mean	Max	Min	Mean	Max		Mean		
Midlatitudes	0.7	3.4	10	0.7	3.4	10	0.7	3.4	10		3.4		
Tropics	2.8	10.2	29.7	2.8	10.2	29.7	2.8	10.2	29.7		10.2		
$\beta_1$					Mean					0	Min	Mean	Max
$\beta_2$					Mean					0	Min	Mean	Max

the experiment without the plume-in-grid development (BC case) are presented.

#### 4.2.1 Lightning emissions and LNO<sub>x</sub> tracer distributions

Figure 3 displays the geographical distributions of the 9 km lightning NO<sub>x</sub> emissions (a), the related LNO<sub>x</sub> tracer distributions (b), and the LNO<sub>x</sub> tracer zonal average (c) in January (top panels) and in July (bottom panels), reproduced by the CTM from the P1 experiment. These results are shown for an approximate detrainment level (9 km altitude) where the detrainment of LNO<sub>x</sub> is the largest. In January, the highest emissions of NO<sub>x</sub> from lightning ( $4\text{--}6 \times 10^9$  molecules cm<sup>-2</sup> s<sup>-1</sup>) are located in the Southern Hemisphere around the tropics over west Australia and central-southern Africa. Also, the model gives low LNO<sub>x</sub> emissions ( $< 3 \times 10^9$  molecules cm<sup>-2</sup> s<sup>-1</sup>) over South America and North America, especially over the Gulf of Mexico. In July, the highest LNO<sub>x</sub> emissions ( $4\text{--}6 \times 10^9$  molecules cm<sup>-2</sup> s<sup>-1</sup>) are calculated in the Northern Hemisphere over North America, the north of India, central Africa, and the Sahel. In addition, LNO<sub>x</sub> emissions are modeled over Europe and over east Asia, but to a lesser extent ( $< 2 \times 10^9$  molecules cm<sup>-2</sup> s<sup>-1</sup>).

The lightning NO<sub>x</sub> tracer introduced into the model represents the lightning NO<sub>x</sub> emissions affected by the transport and the exponential decay depending on the plume lifetime. Figure 3 shows that the tracer distribution is consistent with the lightning NO<sub>x</sub> emissions. However, it is important to note that the plume lifetime is a key factor in the evolution of the LNO<sub>x</sub> tracer mixing ratio. A long plume lifetime (several hours to several days) allows the intercontinental transport of LNO<sub>x</sub> plumes. The representation of the sub-grid chemistry and the transport of the nonlinear chemistry effects related to the plume consideration becomes important for the chemistry of the regions located far downwind from source regions. The plume lifetime depends on the latitude because of the different background chemical concentrations and the different amount of NO<sub>x</sub> emitted from lightning in the tropics and in the midlatitudes. In addition, as mentioned before,

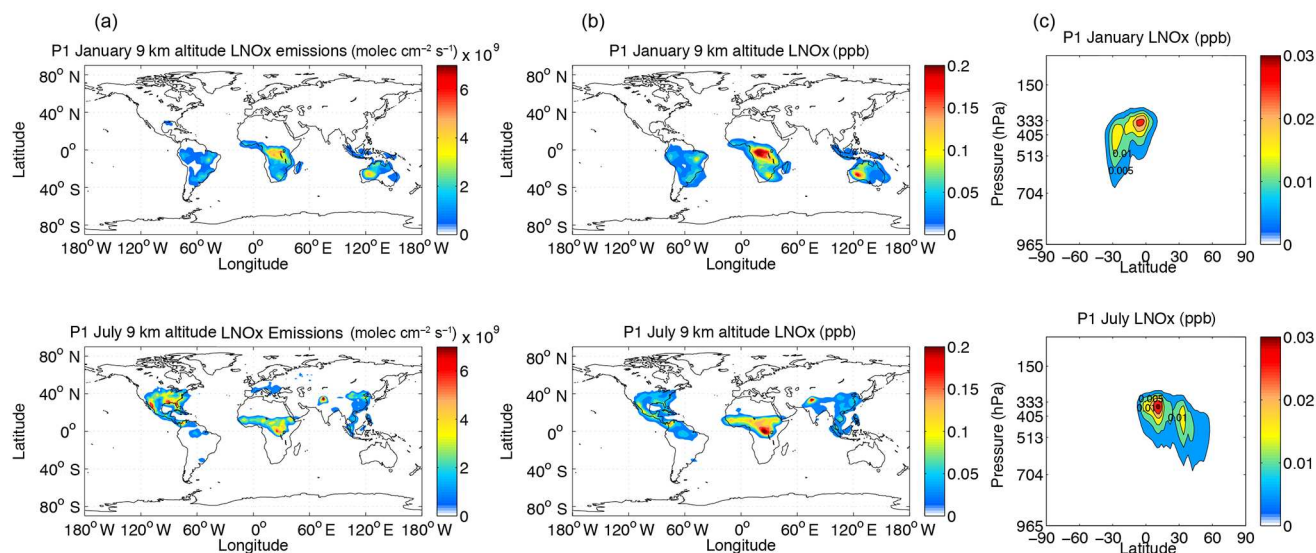
we consider the influence of day and night conditions on the plume lifetime estimate.

According to its preliminary calculation (Sect. 3.2.3), the plume lifetime is longer in the tropics (9 and 21.3 h for day and night conditions, respectively) than in the midlatitudes (3 and 6 h, for day and night conditions, respectively). So, the LNO<sub>x</sub> tracer is characterized by a shorter lifetime as a plume over North America than over central Africa and around the Sahel, while the model simulated fewer emissions over these regions, especially in summer. In boreal winter, the mixing ratio of the lightning NO<sub>x</sub> tracer calculated by the model is about 0.21 ppb over central and southern Africa, 0.18 ppb over west Australia and 0.11 ppb over South America. In summer, the tracer mixing ratio is simulated as 0.21, 0.32, and 0.16 ppb over central Africa, northern India, and North America, respectively. The lightning NO<sub>x</sub> tracer is produced at altitudes where lightning NO<sub>x</sub> is calculated and detrained (in the upper troposphere between  $\sim 500$  and 300 hPa) as shown in Fig. 3c.

#### 4.2.2 Impact of lightning on NO<sub>x</sub> and O<sub>3</sub> distributions with the plume parameterization

The difference between the P1 and BC experiments (P3) was calculated in order to quantify the changes in NO<sub>x</sub> and O<sub>3</sub> mixing ratios on the large scale implied by the implementation of the plume-in-grid parameterization in GEOS-Chem. Figures 4 and 5 display the geographical distributions of the NO<sub>x</sub>, HNO<sub>3</sub>, PAN, and O<sub>3</sub> absolute changes (in ppb) in January and in July, respectively. The 9 km altitude level was chosen because of the most significant variations at this altitude compared to the rest of the troposphere.

In boreal winter, LNO<sub>x</sub> plume chemistry leads to a maximum decrease on a large scale over regions of emissions of 120 ppt for NO<sub>x</sub> and a decrease of 68 ppt for HNO<sub>3</sub> and 16 ppt for PAN over central and southern Africa. These variations are associated with a maximum O<sub>3</sub> decrease of 2.8 ppb over regions of emissions. A similar NO<sub>x</sub>, HNO<sub>3</sub>, PAN, and O<sub>3</sub> reduction is obtained in other areas of high LNO<sub>x</sub> emissions (i.e., over west Australia and South America). Downwind of LNO<sub>x</sub> emissions, the opposite effect is observed for NO<sub>x</sub> and HNO<sub>3</sub> species, with a maximum increase of 40 ppt



**Figure 3.** Geographical distributions at 9 km altitude of lightning NO<sub>x</sub> emissions (left column, **a**), the geographical distributions of the related LNO<sub>x</sub> tracer (in ppb; middle column, **b**), and the zonal average of the LNO<sub>x</sub> tracer (in ppb; right column, **c**) for January (top) and July (bottom). Experiment P1, using  $\tau$  and  $K_{\text{eff}}$  determined with  $D_h = 15 \text{ m}^2 \text{ s}^{-1}$  and  $\text{NO}_i^{\text{mean}}$  and performed with the GEOS-Chem model.

for NO<sub>x</sub> and 13.5 ppt for HNO<sub>3</sub> observed over the South Atlantic and Indian Ocean. Generally, PAN still decreases over oceans but to a lesser extent compared to regions of LNO<sub>x</sub> emissions, with a maximum reduction of 9 ppt. The O<sub>3</sub> response is a maximum increase of 1.13 ppb around the area where the transport is effective and especially over the oceans. In summer, maximum decreases of 140 ppt for NO<sub>x</sub>, 60 ppt for HNO<sub>3</sub>, and 24 ppt for PAN are calculated by the CTM, leading to a maximum O<sub>3</sub> decrease of 2.4 ppb over central Africa (reduction is also observed over North America and northern India). Downwind of lightning emissions, an increase in NO<sub>x</sub> and HNO<sub>3</sub> is observed, with a maximum value of 30 ppt and 38 ppt, respectively. The PAN reservoir species also still decrease slightly downwind, with 2 ppt changes. Finally, this leads to a maximum O<sub>3</sub> increase of 0.7 ppb.

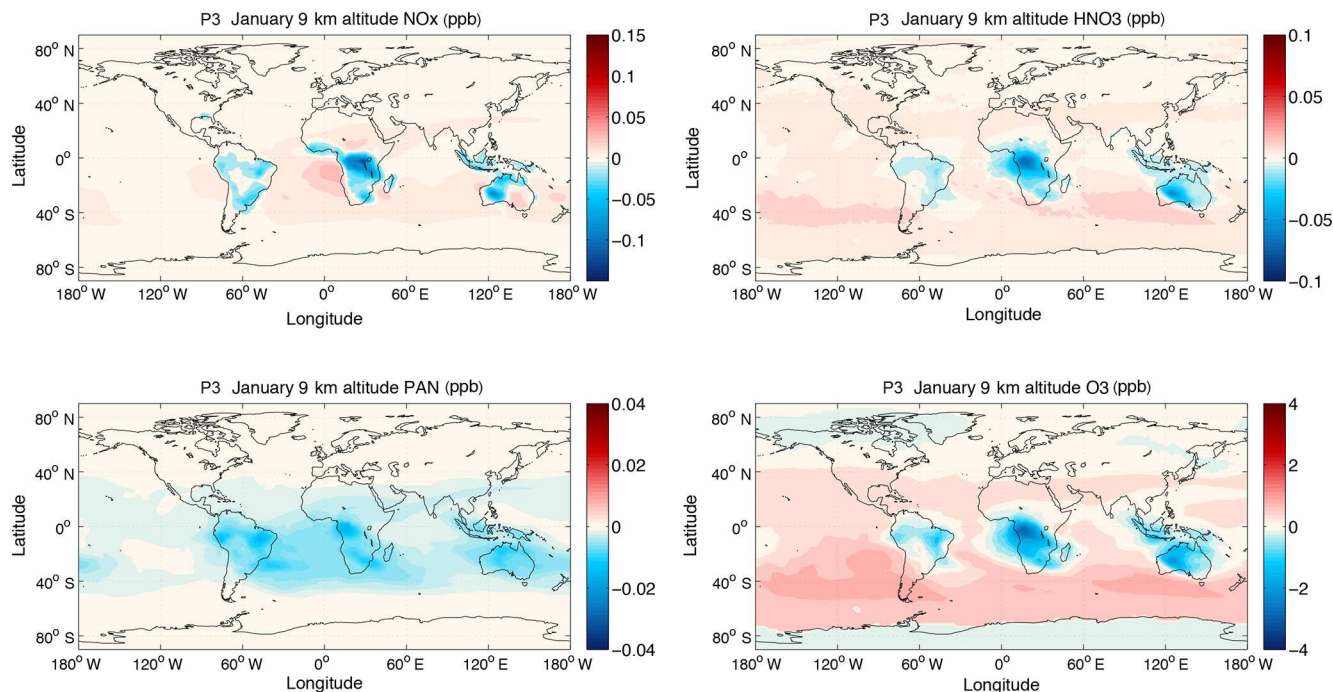
Note that the production of PAN is limited by the supply of NO<sub>x</sub> or non-methane volatile organic compounds (NMVOCs). Above continental lightning sources regions, NMVOCs are uplifted by deep convection but with lower NO<sub>x</sub> due to the activation of the plume parameterization. This implies a less efficient PAN production in these regions. Downwind of lightning source regions (oceanic regions), NO<sub>x</sub> increases because of the LNO<sub>x</sub> transport in the plume, but there are less NMVOCs available to produce PAN. Therefore, both in regions of LNO<sub>x</sub> emissions and downwind, PAN production is limited, leading to overall lower PAN mixing ratios on a large scale in the P1 experiment. However, this a more nuanced view of this may be obtained by considering the PAN chemistry in future studies using a similar LNO<sub>x</sub> plume parameterization and by intro-

ducing the PAN and CH<sub>3</sub>C(O)OO continuity equations and a new term to consider the fraction of NO<sub>x</sub> converted to PAN within the plume. This should enable PAN production during plume transport, which is inhibited in the current version.

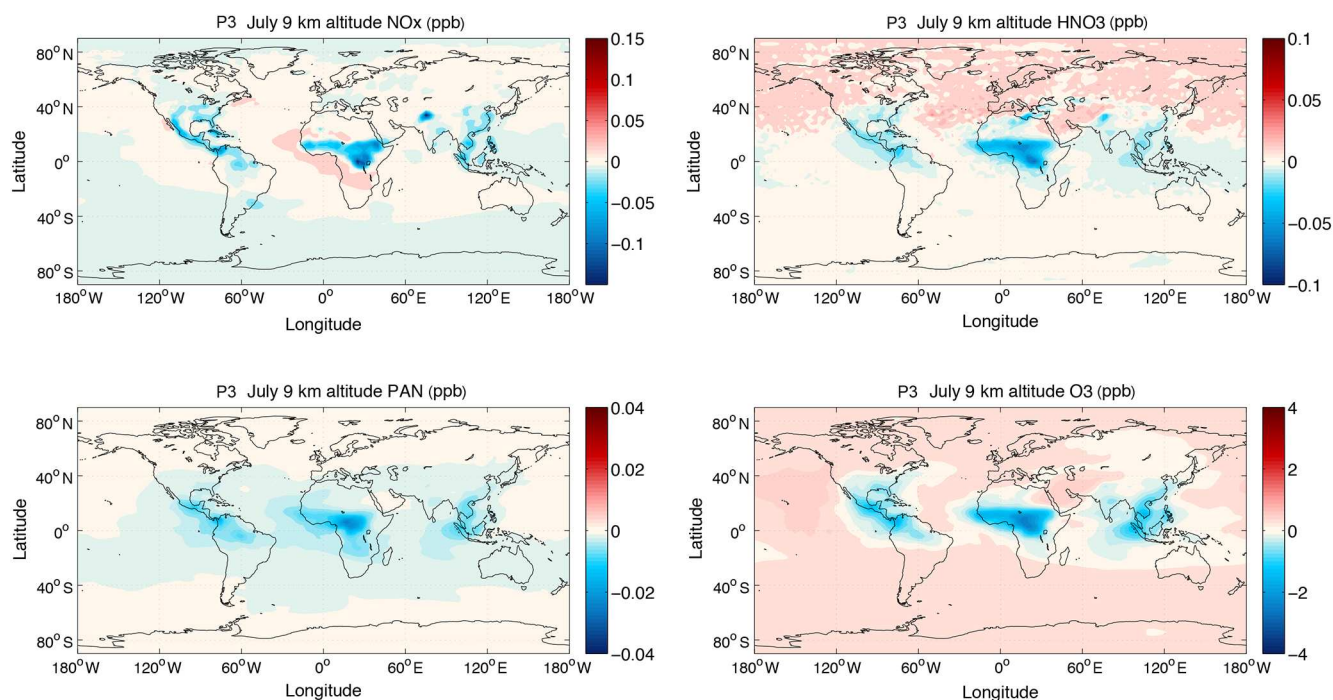
In order to provide a full overview of the effects of the plume parameterization, the relative difference between the P1 and BC experiments (i.e., P3/BC) was calculated integrated throughout the troposphere. Figures 6 and 7 show the zonal average of NO<sub>x</sub> (upper panels) and O<sub>3</sub> (bottom panels) relative changes (in %) integrated throughout the troposphere for the regions of interest for January and July. During boreal winter, the highest NO<sub>x</sub> (O<sub>3</sub>) decreases of 10 % (5 %) in west Australia; then 20 % (6 %) in central Africa are calculated. These negative variations are mainly calculated between 400 hPa and the tropopause level for NO<sub>x</sub> and ozone. South America is characterized by a decrease of 20 % in the nitrogen oxides and 1 % in ozone. Over this region, variations are significant in the entire troposphere for both species. In contrast to the continent decrease, an NO<sub>x</sub> increase is observed over the major part of the South Atlantic and the Indian Ocean, with a 14 and 20 % maximum, respectively. O<sub>3</sub> responds with an increase of 1 % near the tropopause, and it becomes higher by about 4 % close to the surface. In summer, there is an NO<sub>x</sub> (O<sub>3</sub>) decrease of 25 % (8 %) over central Africa, 20 % (2 %) over northern India, and 5 % (0.5 %) over North America. Also, the South Atlantic and Indian Ocean (located downwind of lightning NO<sub>x</sub> emissions) are characterized by a maximum increase of 18 % for NO<sub>x</sub> and 2 % for O<sub>3</sub>.

As a result, the sub-grid chemistry associated with the LNO<sub>x</sub> emissions implies (i) a decrease in the nitrogen oxides



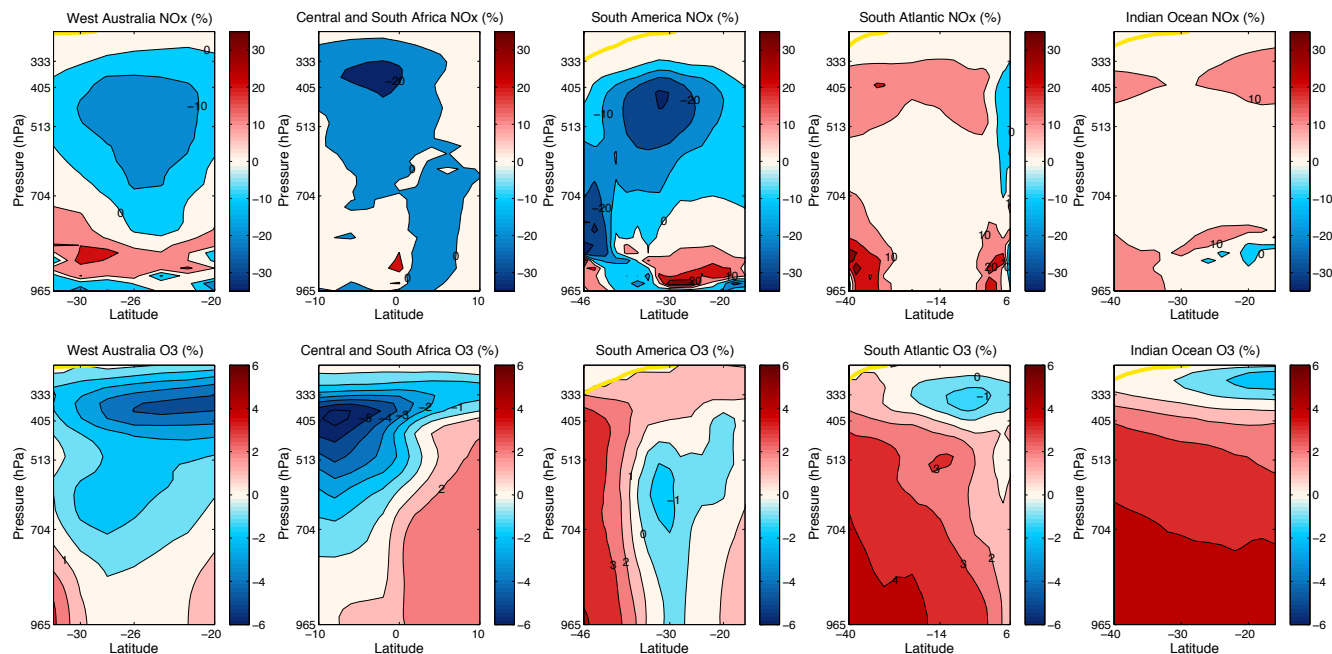


**Figure 4.** Geographical distributions of  $\text{NO}_x$ ,  $\text{HNO}_3$ , PAN, and  $\text{O}_3$  variations (in ppb) at 9 km altitude for January from the absolute difference (P3) between P1 and BC experiments. P1 was performed using  $\tau$  and  $K_{\text{eff}}$  determined with  $D_h = 15 \text{ m}^2 \text{ s}^{-1}$  and  $\text{NO}_i^{\text{mean}}$  with GEOS-Chem.

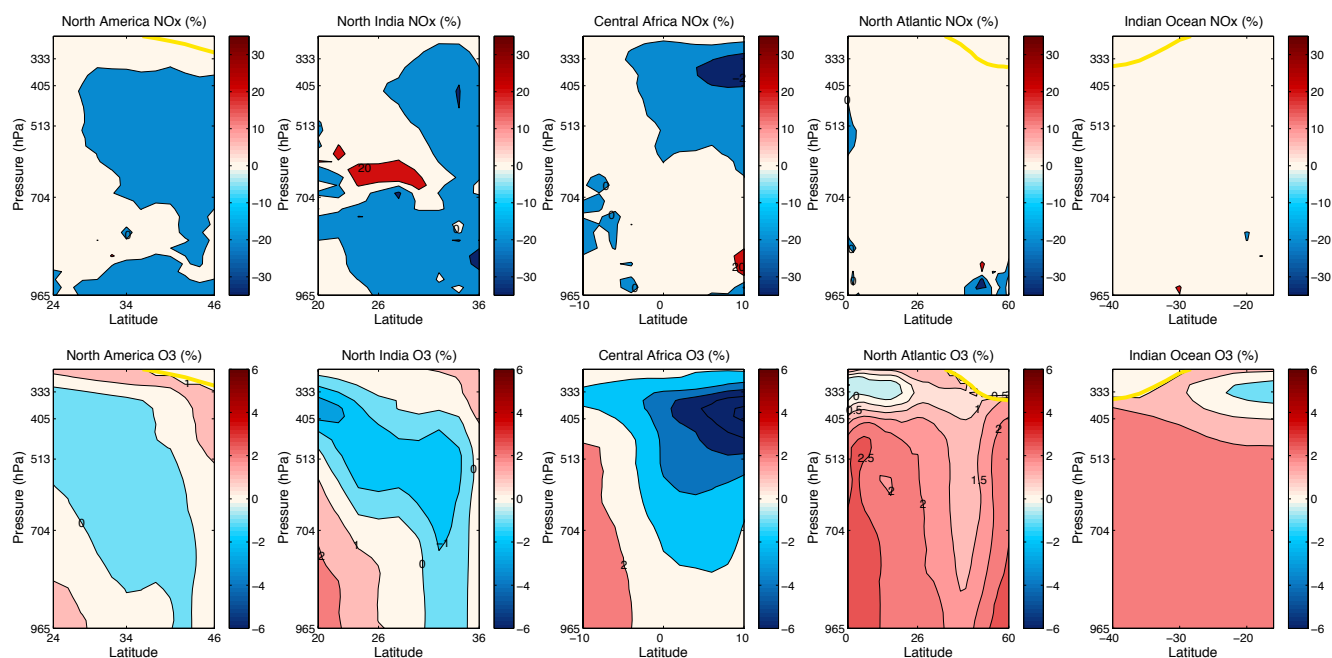


**Figure 5.** Geographical distributions of  $\text{NO}_x$ ,  $\text{HNO}_3$ , PAN, and  $\text{O}_3$  variations (in ppb) at 9 km altitude for July from the absolute difference (P3) between P1 and BC experiments. P1 was performed using  $\tau$  and  $K_{\text{eff}}$  determined with  $D_h = 15 \text{ m}^2 \text{ s}^{-1}$  and  $\text{NO}_i^{\text{mean}}$  with GEOS-Chem.

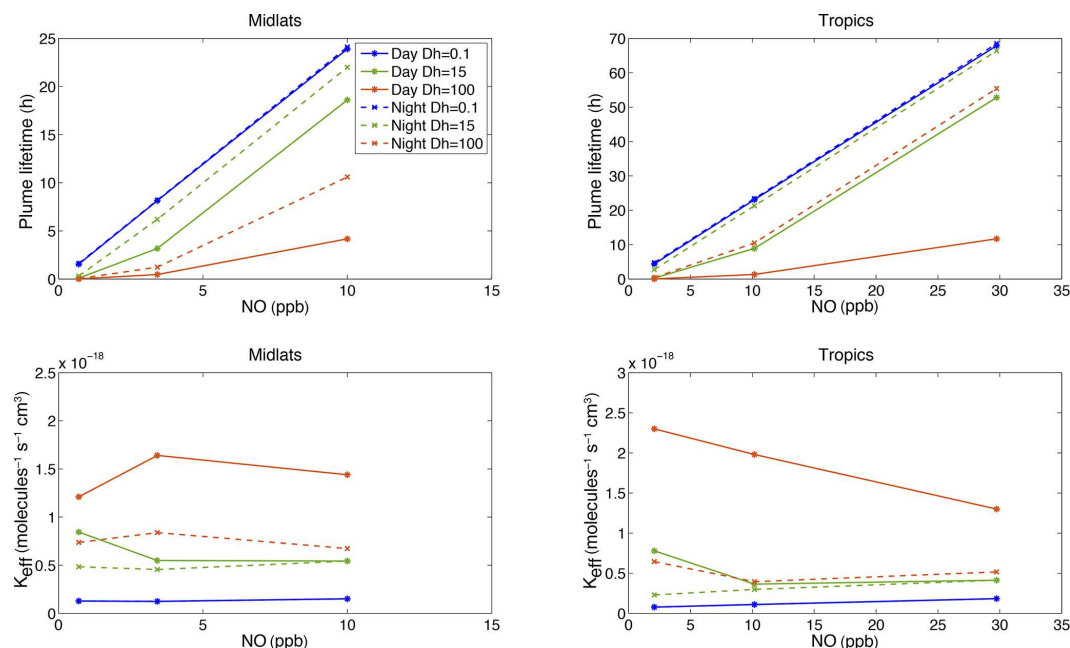




**Figure 6.** Zonal averaged NO<sub>x</sub> (upper panels) and O<sub>3</sub> (bottom panels) variations (in %) over the regions characterized by strong NO<sub>x</sub> emissions for January (the yellow solid line represents the tropopause level) from the relative difference between P1 and BC experiments (P3/BC). P1 was performed using  $\tau$  and  $K_{\text{eff}}$  determined with  $D_h = 15 \text{ m}^2 \text{ s}^{-1}$  and  $\text{NO}_i^{\text{mean}}$  with GEOS-Chem.



**Figure 7.** Zonal averaged NO<sub>x</sub> (upper panels) and O<sub>3</sub> (bottom panels) variations (in %) over the regions characterized by strong NO<sub>x</sub> emissions for July (the yellow solid line represents the tropopause level) from the relative difference between P1 and BC experiments (P3/BC). P1 was performed using  $\tau$  and  $K_{\text{eff}}$  determined with  $D_h = 15 \text{ m}^2 \text{ s}^{-1}$  and  $\text{NO}_i^{\text{mean}}$  with GEOS-Chem.



**Figure 8.** The plume lifetime ( $\tau$ , upper panels) and the effective reaction rate constant ( $K_{\text{eff}}$ , bottom panels) depending (i) on the horizontal coefficient diffusion ( $D_h$ , m<sup>2</sup> s<sup>-1</sup>) for the midlatitudes (left panels) and the tropics (right panels) and (ii) on the NO mixing ratio injected by lightning ( $\text{NO}_i$ , in ppb).

and ozone mixing ratios on a large scale over regions characterized by intense lightning emissions and (ii) an increase in these species downwind of emissions. In particular, the plume parameterization related to the lightning NO<sub>x</sub> leads to

1. significant effects on the NO<sub>x</sub> mixing ratio ( $\pm 20\%$ ): these effects on nitrogen oxides are important because NO<sub>x</sub> is the first criterion which is constrained in a CTM in order to determine the global LNO<sub>x</sub> production (6 TgN yr<sup>-1</sup> in the GEOS-Chem model);
2. lower effects on the O<sub>3</sub> mixing ratio ( $\pm 5\%$ ): these limited impacts on ozone may be explained by the compensatory effect of the NO<sub>y</sub> species (mainly conversion of NO<sub>x</sub> into HNO<sub>3</sub> within the plume).

The effects of the plume parameterization are simulated over the entire troposphere mainly for ozone. Indeed, the spreading of effects on ozone to the lower free troposphere is related to the subsidence areas of the Walker circulation. These regions are characterized by the accumulation and creation of ozone for low-altitude levels. Nevertheless, the maximum NO<sub>x</sub> and O<sub>3</sub> variations are calculated for altitude levels associated with a mean detrainment level. The more realistic representation of the sub-grid processes (P1 experiment) related to the LNO<sub>x</sub> plume is in contrast with the simplified instantaneous dilution in the grid cell of the lightning NO<sub>x</sub> emissions (BC experiment).

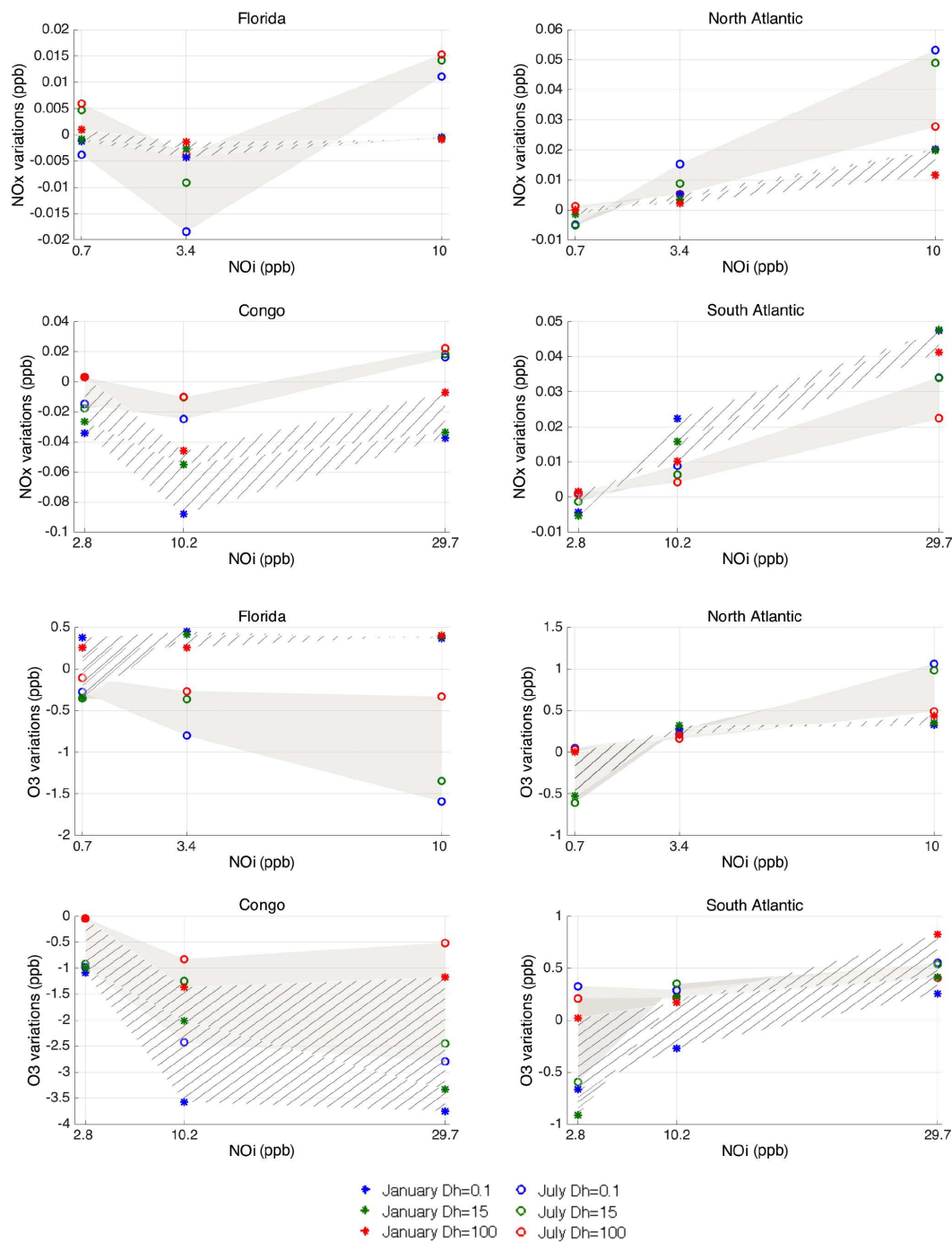
The plume approach allows the conversion of NO<sub>x</sub> into HNO<sub>3</sub> during the plume lifetime. In addition, the high NO<sub>x</sub>

concentration within the plume (much higher than the background content) leads to the O<sub>3</sub> titration and more generally to the O<sub>x</sub> destruction within the plume. The most important impact of the plume parameterization is the transport of the LNO<sub>x</sub> emissions as a plume and the transport of the associated nonlinear chemistry effects, leading to a delay in the O<sub>3</sub> production on a large scale. In other words, less O<sub>3</sub> is produced by photochemical reactions from NO<sub>x</sub> over the regions with intense lightning NO<sub>x</sub> emissions than downwind of LNO<sub>x</sub> emissions.

### 4.3 Plume sensitivity to the estimated uncertainties of parameter calculations

#### 4.3.1 The atmospheric dynamical conditions and the initial NO mixing ratio injected by lightning

The impact of (i) the diffusion properties of the atmosphere ( $D_h$ ) and (ii) the initial NO mixing ratio injected by lightning ( $\text{NO}_i$ ) is analyzed.  $D_h$  and  $\text{NO}_i$  are the two key parameters in the determination of the physical and chemical characteristics of the plume. The modified simulation characterizing the P1 experiment was run for the ranges of the horizontal diffusion coefficients and the initial NO mixing ratio injected by lightning. It is important to note that for these sensitivity tests, the  $\beta_1$  and  $\beta_2$  coefficients remain constant in their mean values. The  $\tau$  and  $K_{\text{eff}}$  values related to these simulations are those previously calculated (Sect. 3) and summarized in Tables 2 and 3. Figure 8 displays  $\tau$  (upper panels) and  $K_{\text{eff}}$



**Figure 9.** The NO<sub>x</sub> and O<sub>3</sub> sensitivity at 9 km altitude depending on the horizontal coefficient diffusion ( $D_h$ ,  $\text{m}^2 \text{s}^{-1}$ ) and on the NO mixing ratio injected by lightning ( $\text{NO}_i$ , ppb) for the midlatitudes (Florida and North Atlantic) and the tropics (Congo Basin and South Atlantic). Intervals are hatched in January and filled in grey in July. Markers correspond to the NO<sub>x</sub> variations simulated for  $D_h = 0.1 \text{ m}^2 \text{s}^{-1}$  (red),  $D_h = 15 \text{ m}^2 \text{s}^{-1}$  (blue), and  $D_h = 100 \text{ m}^2 \text{s}^{-1}$  (green).

**Table 6.** The sensitivity of NO<sub>x</sub> (in ppt) and O<sub>3</sub> (in ppb) depending on the horizontal diffusion coefficient ( $D_h$ , m<sup>2</sup> s<sup>-1</sup>) and on the NO<sub>i</sub> mixing ratio (ppb) injected by lightning for the midlatitudes (Florida and North Atlantic) and tropics (Congo Basin and South Atlantic) in January and July.

	Jan				Jul			
	Midlatitudes		Tropics		Midlatitudes		Tropics	
	Florida	North Atlantic	Congo Basin	South Atlantic	Florida	North Atlantic	Congo Basin	South Atlantic
$\Delta\text{NO}_x \pm$	[-1.7, +1.8]	[-8.2, +1.7]	[-33.1 + 29.7]	[-6.5, +6.9]	[-9.3, +5.4]	[-21.1, +6.6]	[-14.3, +21]	[-11.5, +2.6]
$\Delta\text{O}_3 \pm$	[-0.16, +0.72]	[-0.12, +0.53]	[-1.56, +2.16]	[-0.49, +0.94]	[-0.44, +1.01]	[-0.49, +0.66]	[-1.18, +1.93]	[-0.14, +0.92]

**Table 7.** The sensitivity of NO<sub>x</sub> (in ppt) and O<sub>3</sub> (in ppb) depending on  $\beta_1$  and  $\beta_2$  values for the midlatitudes (Florida and North Atlantic) and tropics (Congo Basin and South Atlantic) in January and July. Experiment P1, using  $D_h = 15 \text{ m}^2 \text{ s}^{-1}$  and NO<sub>i</sub><sup>mean</sup>, performed with the GEOS-Chem model.

	Jan				Jul			
	Midlatitudes		Tropics		Midlatitudes		Tropics	
	Florida	North Atlantic	Congo Basin	South Atlantic	Florida	North Atlantic	Congo Basin	South Atlantic
$\Delta\text{NO}_x \pm \times 10^{-2}$	[-1.6, -0.06]	[-2.3, -2.3]	[-2.3, +0.9]	[+0.3, +0.6]	[-3.3, +1.4]	[-21.1, +6.6]	[+0.4, +2.1]	[-0.9, -0.4]
$\Delta\text{O}_3 \pm \times 10^{-4}$	[-9, +5]	[-9, +4]	[-3, +22]	[-10, +11]	[-24, -6]	[-6, +19]	[-8, +17]	[-30, -2]

(bottom panels) variations depending on  $D_h$  and NO<sub>i</sub>. As expected, the strongest horizontal diffusion leads to the most efficient dispersion of the plume. In both the midlatitudes and the tropics,  $\tau$  decreases when  $D_h$  becomes larger. In addition,  $\tau$  increases with the initial NO mixing ratio injected by lightning. In contrast,  $K_{\text{eff}}$  increases with the  $D_h$  coefficient in the two regions of the globe.

The sensitivity of the NO<sub>x</sub> and O<sub>3</sub> mixing ratios around the mean value for regions and seasons depending on the known uncertainties associated with parameter calculations has been quantified. Figure 9 shows the ranges of sensitivity of NO<sub>x</sub> and O<sub>3</sub> ( $\Delta\text{NO}_x$  and  $\Delta\text{O}_3$ , respectively) at 9 km altitude reproduced by GEOS-Chem depending on  $D_h$  and on the initial NO mixing ratio (NO<sub>i</sub>). Note that for the sake of readability, the scale of NO<sub>x</sub> and O<sub>3</sub> changes differs by region. Results are also summarized in Table 6.

We chose representative continental areas such as Florida and the Congo Basin, which correspond to regions characterized by intense electrical activity for the midlatitudes and the tropics, respectively. The North and South Atlantic were selected to represent regions downwind of NO<sub>x</sub> emissions for the midlatitude and the tropics variations, respectively. The highest NO<sub>x</sub> and O<sub>3</sub> ranges are obtained for continental tropical regions with  $\Delta\text{NO}_x$  [-33.1, +29.7] ppt and  $\Delta\text{O}_3$  [-1.56, +2.16] ppb, in January, and  $\Delta\text{NO}_x$  [-14.3, +21] ppt and  $\Delta\text{O}_3$  [-1.18, +1.93] ppb, in July. The largest range associated with the tropical continents may be explained by the largest difference in parameter values defining the plume in this region (especially NO<sub>i</sub>). The smallest changes are observed over continental midlatitude regions for winter with  $\Delta\text{NO}_x$  [-1.7, +1.8] ppt and  $\Delta\text{O}_3$  [-0.16, +0.72] ppb and over oceanic tropical regions in summer such that  $\Delta\text{NO}_x$  [-11.5, +2.6] ppt and  $\Delta\text{O}_3$  [-0.14, +0.92] ppb. As a result, the sensitivity of NO<sub>x</sub> and O<sub>3</sub> species to the pa-

rameter uncertainties is a few parts per trillion for NO<sub>x</sub> and less than 2 ppb for O<sub>3</sub>.

#### 4.3.2 Coefficients related to the nitrification mechanism ( $\beta_1$ and $\beta_2$ )

In order to estimate the sensitivity of the NO<sub>x</sub> and O<sub>3</sub> mixing ratios related to the uncertainties in the  $\beta_1$  and  $\beta_2$  fractions (Table 7), the difference between the P1 experiment using the  $\beta_1$  and  $\beta_2$  mean values and the P1 experiment using minimum and maximum  $\beta_1$  and  $\beta_2$  coefficients has been calculated. This implies that  $\tau$  and  $K_{\text{eff}}$  are constant.

In January, the highest sensitivity in NO<sub>x</sub> mixing ratio is  $\Delta\text{NO}_x$  [-2.3, +0.9]  $\times 10^{-2}$  ppt over the continental tropical regions and  $\Delta\text{O}_3$  [-10, +11]  $\times 10^{-4}$  ppb over the tropical ocean in O<sub>3</sub>, while the midlatitude oceanic areas show minimum ranges in NO<sub>x</sub> and O<sub>3</sub>, with  $\Delta\text{NO}_x \pm 2.3 \times 10^{-2}$  ppt associated with  $\Delta\text{O}_3$  [-9, +4]  $\times 10^{-4}$  ppb. In July, the maximum ranges are calculated over oceans in the midlatitudes for NO<sub>x</sub>, such that  $\Delta\text{NO}_x$  [-21.1, +6.6]  $\times 10^{-2}$  ppt, and in the tropics for O<sub>3</sub>, with  $\Delta\text{O}_3$  [-30, -2]  $\times 10^{-4}$  ppb. Finally, the smallest changes,  $\Delta\text{NO}_x$  [-0.9, -0.4]  $\times 10^{-2}$  ppt and  $\Delta\text{O}_3$  [-24, -6]  $\times 10^{-4}$  ppb, are simulated for the tropical ocean and the midlatitude continents, respectively.

In addition, the impact of the nitrification mechanism was assessed comparing the P1 experiment using mean  $\beta_1$  and  $\beta_2$  values and the P2 experiment for which  $\beta_1 = \beta_2 = 0$ . As a result, taking into account NO<sub>x</sub> conversion into HNO<sub>3</sub> and using the mean  $\beta$  fractions calculated in this study does not imply strong changes in NO<sub>x</sub> and O<sub>3</sub> distributions ( $\Delta\text{NO}_x < 10^{-4}$  ppt and  $\Delta\text{O}_3 < 10^{-2}$  ppb).

In the case of significant values of the  $\beta$  fractions, the rate of the nitrification mechanism should imply a delay in the O<sub>3</sub> formation from the NO<sub>x</sub> in the plume because of the NO<sub>x</sub> storage in HNO<sub>3</sub>. On the other hand, HNO<sub>3</sub> is considered one

of the main sinks for NO<sub>x</sub> species undergoing wet deposition and seemingly limiting their affect on global ozone.

The sensitivity tests indicate the limited effect of the NO<sub>x</sub> conversion to HNO<sub>3</sub> within the plume using our  $\beta_1$  and  $\beta_2$  calculations. The sensitivity on NO<sub>x</sub> and O<sub>3</sub> mixing ratios related to  $\beta$  coefficients is about a few parts per trillion. Indeed, our  $\beta_1$  and  $\beta_2$  estimates are smaller than those calculated by Cariolle et al. (2009) ( $\beta_1 = 0.06$  and  $\beta_2 = 0.2$ ), which showed a large impact of this mechanism in the case of aircraft NO<sub>x</sub> emissions. In the present study, we can easily suppose that the increase in the  $\beta_1$  and  $\beta_2$  coefficients should be in agreement with the work of Cooper et al. (2014) in reducing the underestimate of the HNO<sub>3</sub> production induced by NO<sub>x</sub> emissions from lightning. Further estimates of  $\beta$  should be realized using future observations in the cloud anvil of primary species, aerosols, and particles needed for NO<sub>x</sub> conversion during the daytime and nighttime to improve the determination of these parameters. The  $\beta_1$  coefficient is particularly dependent on the HO<sub>x</sub> radicals, which may vary significantly within the cloud anvil, in part because of the transport of peroxides from the lower troposphere by convective uplift (Wennberg et al., 1998). The determination of  $\beta_2$ , corresponding to the NO<sub>x</sub> conversion fraction into HNO<sub>3</sub> via N<sub>2</sub>O<sub>5</sub> formation during nighttime is considerably dependent on (i) the estimate of aerosols and the ice crystal concentration and their lifetime within the cloud anvil, which is highly uncertain according to measurement campaigns, and (ii) the reaction probability regarding aerosols  $\gamma_{\text{N}_2\text{O}_5}^{\text{aerosol}}$  and ice crystals  $\gamma_{\text{N}_2\text{O}_5}^{\text{ice}}$  from laboratory study extrapolations.

According to the results presented in this section, the sensitivity tests show the predominance of the initial NO mixing ratio injected by lightning (NO<sub>i</sub>) and the diffusion properties of the atmosphere ( $D_h$ ) in the variability of the NO<sub>x</sub> and O<sub>3</sub> mixing ratios around the mean value in response to the plume-in-grid parameterization in the CTM. In winter, the NO<sub>x</sub> and O<sub>3</sub> sensitivity is the highest for continental regions in the tropics and the smallest sensitivity is calculated for the midlatitudes. In summer, the most important sensitivity of NO<sub>x</sub> and O<sub>3</sub> is simulated in the tropics over regions characterized by intense LNO<sub>x</sub> emissions, while the least significant sensitivity is still obtained in the tropics but downwind of emissions (mainly over oceans).

## 5 Conclusions

For the first time, a more realistic lightning NO<sub>x</sub> chemistry is implemented as a plume parameterization in a global chemical transport model. The key parameters characterizing the lightning-related plume were estimated depending on two main criteria, i.e., the NO mixing ratio injected by lightning (NO<sub>i</sub>) and the atmospheric diffusion coefficient ( $D_h$ ).

According to the NO<sub>i</sub> and  $D_h$  ranges, the plume lifetime ( $\tau$ ) and the effective reaction rate constant ( $K_{\text{eff}}$ ) for NO<sub>x</sub>–O<sub>3</sub> chemical interactions were estimated as follows:

$$- \tau = [0.01, 68.5] \text{ h};$$

$$- K_{\text{eff}} = [0.77, 23] \times 10^{-19} \text{ molecules}^{-1} \text{ s}^{-1} \text{ cm}^3.$$

Also, for the conditions defined by NO<sub>i</sub><sup>mean</sup> and  $D_h = 15 \text{ m}^2 \text{ s}^{-1}$

$$- \tau \text{ is } 3 \text{ (6) h in the midlatitudes and } 9 \text{ (21.3) h in the tropics during the daytime (nighttime);}$$

$$- K_{\text{eff}} \text{ is } 5.49 \times 10^{-19} \text{ molecules}^{-1} \text{ s}^{-1} \text{ cm}^3 \text{ (} 4.55 \times 10^{-19} \text{ molecules}^{-1} \text{ s}^{-1} \text{ cm}^3 \text{) in the midlatitudes and } 3.64 \times 10^{-19} \text{ molecules}^{-1} \text{ s}^{-1} \text{ cm}^3 \text{ (} 2.98 \times 10^{-19} \text{ molecules}^{-1} \text{ s}^{-1} \text{ cm}^3 \text{) in the tropics during the daytime (nighttime).}$$

Finally, the fractions of NO<sub>x</sub> conversion into HNO<sub>3</sub> within the plume are  $\beta_1 = [1.34, 1.88] \times 10^{-4}$  and  $\beta_2 = [0.24, 14.4] \times 10^{-3}$  for day and night conditions, respectively.

GEOS-Chem simulations performed using mean values for NO<sub>i</sub> and  $D_h = 15 \text{ m}^2 \text{ s}^{-1}$  reveal nitrogen species and ozone changes compared to the instantaneous dilution. A decrease in NO<sub>x</sub> and O<sub>3</sub> mixing ratios on a large scale over regions of strong LNO<sub>x</sub> emissions is observed mainly in the Northern Hemisphere in summer and in the Southern Hemisphere in winter. In the troposphere, a maximum decrease of 20 % (6 %) in January and 25 % (8 %) in July for NO<sub>x</sub> (O<sub>3</sub>) is found over central Africa. In contrast, an increase in NO<sub>x</sub> (O<sub>3</sub>) downwind of emissions of 20 % (4 %) in January and 18 % (2 %) in July is simulated. The LNO<sub>x</sub> plume parameterization allows the transport of the effects on the nonlinear chemistry occurring within the plume and the conversion of NO<sub>x</sub> to nitrogen reservoir species (mainly HNO<sub>3</sub>). However, the most significant impact is the transport of the LNO<sub>x</sub> as a plume. This implies a delay of (i) the NO<sub>x</sub> release into the point grid and (ii) ozone production from NO<sub>x</sub> emitted by lightning flashes corresponding to the decrease in the NO<sub>x</sub> and O<sub>3</sub> mixing ratios on a large scale over regions of emissions and their increase over the transport pathway.

The sensitivity of the NO<sub>x</sub> and O<sub>3</sub> mixing ratios around the mean value depending on the known uncertainties in the plume physics and chemistry key parameters has been estimated. The highest sensitivity is obtained for the continental tropical regions with  $\Delta\text{NO}_x$  [−33.1, +29.7] ppt and  $\Delta\text{O}_3$  [−1.56, +2.16] ppb, in January, and  $\Delta\text{NO}_x$  [−14.3, +21] ppt and  $\Delta\text{O}_3$  [−1.18, +1.93] ppb, in July. Concerning the  $\beta_1$  and  $\beta_2$  fractions, the highest sensitivity depending on the fraction uncertainties for NO<sub>x</sub> is  $\Delta\text{NO}_x$  [−2.3, +0.9]  $\times 10^{-2}$  ppt over the continental tropical regions and  $\Delta\text{O}_3$  [−10, +11]  $\times 10^{-4}$  ppb for O<sub>3</sub> over the tropical ocean in January. In summer, the maximum ranges are calculated over oceans in the midlatitudes for NO<sub>x</sub>, such as  $\Delta\text{NO}_x$  [−21.1, +6.6]  $\times 10^{-2}$  ppt, and in the tropics for O<sub>3</sub>, with  $\Delta\text{O}_3$  [−30, −2]  $\times 10^{-4}$  ppb. Accordingly, the parameters leading to the highest uncertainties in results and those which drive the plume-in-grid parameterization are NO<sub>i</sub> and  $D_h$ .

This study demonstrates the importance of considering the plume-in-grid chemistry related to the lightning NO<sub>x</sub> emissions occurring on a smaller scale for global calculations. Taking into account the plume dilution into the background atmosphere in time and space with the transport of the NO<sub>x</sub> and O<sub>3</sub> nonlinear chemistry effects and the conversion of NO<sub>x</sub> into HNO<sub>3</sub> reservoir species implies more realistic NO<sub>x</sub> and O<sub>3</sub> concentrations in CTM. By allowing a more realistic sub-grid chemistry, the plume-in-grid approach will allow the improvement of the different steps in lightning NO<sub>x</sub> emissions modeling, such as the convection process, the calculation of the NO molecules produced by lightning discharges depending on regions according to recent and future satellite observations, and also processes such as HNO<sub>3</sub> scavenging and HNO<sub>3</sub> uptake by ice crystals.

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